# Construction and reconstruction of tight framelets and wavelets via matrix mask functions ${ }^{\text {st }}$ 

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#### Abstract

The paper develops construction procedures for tight framelets and wavelets using matrix mask functions in the setting of a generalized multiresolution analysis (GMRA). We show the existence of a scaling vector of a GMRA such that its first component exhausts the spectrum of the core space near the origin. The corresponding low-pass matrix mask has an especially advantageous form enabling an effective reconstruction procedure of the original scaling vector. We also prove a generalization of the Unitary Extension Principle for an infinite number of generators. This results in the construction scheme for tight framelets using lowpass and high-pass matrix masks generalizing the classical MRA constructions. We prove that our scheme is flexible enough to reconstruct all possible orthonormal wavelets. As an illustration we exhibit a pathwise connected class of non-MSF non-MRA wavelets sharing the same wavelet dimension function.


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## 1. Introduction and preliminaries

The main aim of this work is to develop a constructive procedure for constructing tight framelets and wavelets from more primitive objects given by low-pass and high-pass matrix mask functions. We should add that all of our results are shown in the general setting of expansive integer dilations in $\mathbb{R}^{n}$. The novelty of our approach lies in its versatility allowing construction of all possible orthonormal wavelets without the customary restrictions on smoothness or decay in time or frequency domains. Hence, it applies to all sorts of exotic and little understood wavelets such as those with unbounded wavelet dimension function. In the case of multiresolution analysis (MRA) wavelets such procedure is well studied and understood.

Usually, an MRA construction starts with a 1-periodic measurable function $m$, also called a low-pass mask and satisfying the quadrature-mirror equation

$$
\begin{equation*}
|m(\xi)|^{2}+|m(\xi+1 / 2)|^{2}=1 \quad \text { for a.e. } \xi \in \mathbb{R} . \tag{1.1}
\end{equation*}
$$

Under small regularity assumptions, such as $m$ is Hölder continuous at 0 and $m(0)=1$, one defines a refinable function $\varphi \in L^{2}(\mathbb{R})$ by

$$
\begin{equation*}
\hat{\varphi}(\xi)=\prod_{j=1}^{\infty} m\left(2^{-j} \xi\right) \tag{1.2}
\end{equation*}
$$

While $\varphi$ might fail to be an orthogonal scaling function of an MRA, one can always obtain a tight frame wavelet $\psi \in L^{2}(\mathbb{R})$ using a high-pass mask $h$ by

$$
\begin{equation*}
\hat{\psi}(\xi)=h(\xi / 2) \hat{\varphi}(\xi / 2), \quad \text { where } h(\xi)=e^{2 \pi i \xi} \overline{m(\xi+1 / 2)}, \tag{1.3}
\end{equation*}
$$

see [20,26]. The fact that $\psi$ is a tight framelet can be shown directly by employing the characterization equations [25, Section 7.1]. Alternatively, it is also a consequence of a Unitary Extension Principle of Ron and Shen [21,29].

While Hölder continuity of $m$ at 0 is a relatively weak assumption, some MRA wavelets cannot be obtained by this scheme [18]. To circumvent this problem, Paluszyński, Šikić, Weiss, and Xiao [27] introduced the class of low-pass filters satisfying

$$
\lim _{n \rightarrow \infty} \prod_{j=n}^{\infty}\left|m\left(2^{-j} \xi\right)\right|=1 \quad \text { for a.e. } \xi \in \mathbb{R}
$$

This is obviously a much weaker condition than Hölder continuity. Moreover, any low-pass mask $m$ of an MRA scaling function must satisfy it by the characterization of scaling functions of MRAs [25]. While the infinite product (1.2) might not be convergent, the authors of [27] proved that one can always construct a refinable function $\varphi$ satisfying

$$
\hat{\varphi}(\xi)=m(\xi / 2) \hat{\varphi}(\xi / 2) \quad \text { a.e. } \xi \in \mathbb{R} .
$$

This is because the product (1.2) converges after taking absolute values and one must only recover the phase factor of $\hat{\varphi}$ by using a multiplier argument. As a consequence, the procedure of constructing MRA tight frame wavelets from [27] recovers all possible MRA wavelets.

Similar ideas were used to prove the connectivity result for MRA wavelets by the Wutam Consortium [34].

Since MRA wavelets form only a special class among all orthonormal wavelets, one could ask whether similar construction and reconstruction procedures are possible for non-MRA wavelets. The most natural way of classifying non-MRA wavelets uses the wavelet dimension function

$$
D_{\psi}(\xi)=\sum_{j=1}^{\infty} \sum_{k \in \mathbb{Z}}\left|\hat{\psi}\left(2^{j}(\xi+k)\right)\right|^{2}
$$

The wavelet dimension function has many interesting properties that were investigated by several authors $[1,2,17,25,28,30]$. One of its fundamental properties says that $D_{\psi}$ can be identified with the multiplicity function of the core space of a GMRA generated by the wavelet $\psi$. In particular, $\psi$ is an MRA wavelet if and only if $D_{\psi} \equiv 1$. Moreover, by the result of Speegle and the authors [17] all possible wavelet dimension functions $\mathcal{D}$ are characterized by the following 4 conditions:
(D1) $\mathcal{D}: \mathbb{R} \rightarrow \mathbb{N} \cup\{0\}$ is a measurable 1-periodic function,
(D2) $\mathcal{D}(\xi)+\mathcal{D}(\xi+1 / 2)=\mathcal{D}(2 \xi)+1$ for a.e. $\xi \in \mathbb{R}$,
(D3) $\sum_{k \in \mathbb{Z}} \mathbf{1}_{\Delta}(\xi+k) \geqslant \mathcal{D}(\xi)$ for a.e. $\xi \in \mathbb{R}$, where

$$
\Delta=\left\{\xi \in \mathbb{R}: \mathcal{D}\left(2^{-j} \xi\right) \geqslant 1 \text { for } j \in \mathbb{N} \cup\{0\}\right\}
$$

(D4) $\liminf _{j \rightarrow \infty} \mathcal{D}\left(2^{-j} \xi\right) \geqslant 1$ for a.e. $\xi \in \mathbb{R}$.
Note that we have intentionally omitted the integrability condition on $\mathcal{D}$, since it is a consequence of (D1) and (D2) by Lemma 3.1 proved in this paper.

The above characterization opens the possibility of constructing wavelets and framelets from more general low-pass matrix masks than the standard scalar masks satisfying (1.1). In general, one starts with a measurable matrix-valued 1-periodic function $M$ satisfying

$$
\begin{equation*}
M(\xi) M^{*}(\xi)+M(\xi+1 / 2) M^{*}(\xi+1 / 2)=\Omega(2 \xi) \quad \text { for a.e. } \xi \in \mathbb{R} \tag{1.4}
\end{equation*}
$$

where

$$
\Omega(\xi)=\operatorname{diag}\left(\mathbf{1}_{S_{1}}(\xi), \mathbf{1}_{S_{2}}(\xi), \ldots\right), \quad S_{j}=\{\xi \in \mathbb{R}: \mathcal{D}(\xi) \geqslant j\}, \quad j \in \mathbb{Z}
$$

More precisely, the values of $M(\xi)$ are infinite size matrices (doubly indexed by $\mathbb{N}$ ) with only finitely many non-zero entries, which can be identified with bounded operators on $\ell^{2}(\mathbb{N})$. Furthermore, in the case when the multiplicity function is bounded by $N \in \mathbb{N}$, we can safely assume that the values of $M(\xi)$ are $N \times N$ matrices.

Baggett, Jorgensen, Merrill, and Packer [4] showed that if the multiplicity function $\mathcal{D}$ is bounded and $M$ satisfies some weak regularity assumptions then one can define a refinable vector function $\Phi=\left(\varphi_{j}\right)_{j \in J} \subset L^{2}(\mathbb{R}), J=\{1, \ldots, N\}$, by

$$
\begin{equation*}
\hat{\Phi}(\xi)=\left[\prod_{j=1}^{\infty} M\left(2^{-j} \xi\right)\right] e \tag{1.5}
\end{equation*}
$$

where $e=(1,0, \ldots, 0)$. To make sure that the above product converges the authors of [4] assume that $M$ is Lipschitz continuous at 0 and $M(0)$ is a matrix having all zero entries except a single 1 in the upper-left corner. In general, $\Phi$ might not be a scaling vector of some GMRA in the sense that each $\varphi_{j}$ is a quasi-orthogonal generator of

$$
\mathcal{S}\left(\varphi_{j}\right)=\overline{\operatorname{span}}\left\{\varphi_{j}(\cdot-k): k \in \mathbb{Z}\right\}
$$

and $\mathcal{S}\left(\varphi_{i}\right) \perp \mathcal{S}\left(\varphi_{j}\right)$ for $i \neq j$. Nevertheless, the authors of [4] proved that by choosing an appropriate high-pass matrix mask $H$, one can always obtain a tight frame wavelet $\psi \in L^{2}(\mathbb{R})$ by setting

$$
\begin{equation*}
\hat{\psi}(\xi)=H(\xi / 2) \hat{\Phi}(\xi / 2) \quad \text { a.e. } \xi \in \mathbb{R} \tag{1.6}
\end{equation*}
$$

More precisely, $H$ is 1-periodic measurable $1 \times N$ matrix-valued function satisfying

$$
\begin{array}{ll}
H(\xi) H^{*}(\xi)+H(\xi+1 / 2) H^{*}(\xi+1 / 2)=1 & \text { a.e. } \xi \in \mathbb{T} \\
M(\xi) H^{*}(\xi)+M(\xi+1 / 2) H^{*}(\xi+1 / 2)=0 & \text { a.e. } \xi \in \mathbb{T} \tag{1.8}
\end{array}
$$

This naturally leads to a fundamental problem of the theory of non-MRA wavelets, which asks whether it is possible to use the above scheme of low-pass and high-pass matrix masks to construct all orthonormal wavelets.

The goal of this paper is to give an affirmative answer to this problem. To give the idea of the level of difficulty behind this project one should realize that, a priori, no regularity assumption on the low-pass matrix mask functions can be assumed. Furthermore, an example in [17] demonstrates that the multiplicity function $\mathcal{D}$ could be unbounded which leads to a matrix-valued low-pass mask $M$ of infinite size. Hence, the infinite product in (1.5) might not converge and special convergence procedures are needed to interpret such ill-defined expressions.

The starting point of this paper is the investigation of the properties of scaling vectors corresponding to the core space of a GMRA. Unlike the case of an MRA, where the scaling function is unique (up to a unimodular 1-periodic phase factor in the Fourier domain), there are many possible choices for scaling vectors for a GMRA. This has been traditionally considered as an impediment of a successful theory, since different choices of a scaling vector $\Phi$ could lead to totally different low-pass masks $M$ satisfying

$$
\begin{equation*}
\hat{\Phi}(\xi)=M(\xi / 2) \Phi(\xi / 2) \quad \text { a.e. } \xi \in \mathbb{R} . \tag{1.9}
\end{equation*}
$$

It might seem that the only useful information extracted from (1.9) is a matrix analogue of the quadrature-mirror equation (1.4). Nevertheless, we show that abundance of choice is also a blessing if one carefully chooses generators of the scaling vector. The key idea is to choose the first generator $\varphi_{1}$ such that it exhausts the entire spectrum of the core space near the origin. Consequently, the remaining generators $\varphi_{2}, \varphi_{3}, \ldots$ must be supported away from the origin in the Fourier domain. This leads to an especially advantageous form of the low-pass matrix mask $M$ such that the first column of $M(\xi)$ has zeros in every entry, except the first where it has absolute value "almost equal" to 1 for $\xi \approx 0$. For a precise statement see Theorem 2.2.

In the case when a GMRA is associated with an orthonormal wavelet $\psi$, it is easy to verify that the high-pass mask $H$ given by (1.6) must satisfy (1.7) and (1.8). The crux of our approach
is the assertion claiming that one can reverse the above process. That is, given a low-pass mask in the above advantageous form and a high-pass matrix mask satisfying (1.7) and (1.8), we can construct an associated tight frame wavelet $\psi$. This leads to Theorem 4.3 which is the main construction result of our paper. The first key ingredient in the proof of this result is the existence of a refinable vector $\Phi$, which is a result of a special convergence procedure making sense out of potentially divergent infinite product in (1.5). The second ingredient is a generalization of the Unitary Extension Principle to a situation when a refinable vector $\Phi$ has infinitely many components.

The last part of the paper proves that the above scheme is flexible enough to reconstruct all possible wavelets $\psi$. A pivotal role in the reconstruction scheme is played by the concept of a multiplier. We say that a unimodular function $v$ is a multiplier associated to $M=\left(m_{i, j}\right)$ if it satisfies

$$
\nu(2 \xi) \overline{\nu(\xi)}\left|m_{1,1,}(\xi)\right|=m_{1,1}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}
$$

Then, our main reconstruction Theorem 5.4 says that the scaling vector $\Phi$ can be recovered by

$$
\begin{equation*}
\hat{\Phi}(\xi)=\lim _{N \rightarrow \infty} v\left(2^{-N} \xi\right)\left[\prod_{j=1}^{N} M\left(2^{-j} \xi\right)\right] e \quad \text { for a.e. } \xi \tag{1.10}
\end{equation*}
$$

and then the wavelet $\psi$ can be recovered by (1.6).
Finally, the paper ends with examples illustrating the inner workings of our construction and reconstruction procedures. In particular, we give an example of a class $\mathcal{W}_{\text {nik }}$ of non-MSF and non-MRA wavelets such that all of its members share the same Journé dimension function. Recall that a wavelet $\psi$ is said to be minimally supported frequency (MSF), if the Lebesgue measure of the support of $\hat{\psi}$ is smallest possible, that is equal to 1 . The classes of MSF wavelets and MRA wavelets are already well studied and understood. However, the class of non-MSF and non-MRA wavelets is the least understood and has inhibited the growth of $L^{2}$ theory of wavelets. Nevertheless, we prove that our class $\mathcal{W}_{\text {nik }}$ is pathwise connected indicating that our techniques have a potential of attacking a recalcitrant problem of the connectivity of the set of all orthonormal wavelets.

Despite the fact that all of our results are motivated by the classical case of dyadic dilations in $\mathbb{R}$, we will adopt a more general setting of expansive integer-valued dilations in $\mathbb{R}^{n}$. More specifically, we shall assume that we are given an $n \times n$ integer-valued matrix $A$ that is expansive, i.e., all its eigenvalues have modulus greater than 1 . For simplicity, its transpose will be denoted by $B$.

We recall that a sequence $\left\{D^{j}(V): j \in \mathbb{Z}\right\}$ of closed subspaces of $L^{2}\left(\mathbb{R}^{n}\right)$ is called a generalized multiresolution analysis (GMRA) if
(M1) $T_{k} V=V$ for all $k \in \mathbb{Z}^{n}$,
(M2) $V \subset D(V)$,
(M3) $\overline{\bigcup_{j \in \mathbb{Z}} D^{j}(V)}=L^{2}\left(\mathbb{R}^{n}\right)$,
(M4) $\bigcap_{j \in \mathbb{Z}} D^{j}(V)=\{0\}$.
Here, the dilation operator $D$ is given by $D f(x)=|\operatorname{det} A|^{1 / 2} f(A x)$ for some $n \times n$ expansive integer-valued matrix $A$ and the translation operator $T_{k} f(x)=f(x-k)$ for some $k \in \mathbb{Z}^{n}$.

As we can see, a GMRA is based on the core space $V$. Condition (M1) means that $V$ is a shift-invariant (SI) space. If $V$ satisfies (M2), then we call it refinable. Also, we shall often write $V_{j}$ instead of $D^{j}(V)$.

We say that a finite family $\Psi=\left\{\psi^{1}, \ldots, \psi^{N}\right\} \subset L^{2}\left(\mathbb{R}^{n}\right)$ is a wavelet if its associated affine system

$$
\psi_{j, k}(x)=|\operatorname{det} A|^{j / 2} \psi\left(A^{j} x-k\right), \quad j \in \mathbb{Z}, k \in \mathbb{Z}^{n}, \psi \in \Psi
$$

is an orthonormal basis of $L^{2}\left(\mathbb{R}^{n}\right)$. In the more general case, when the affine system is a tight frame (with constant 1), we say that $\Psi$ is a tight framelet. The latter are characterized by the wellknown equations that we list in (3.7) and (3.8). Moreover, a framelet $\Psi$ is called semi-orthogonal if

$$
\bigoplus_{j \in \mathbb{Z}} D^{j}(W)=L^{2}\left(\mathbb{R}^{n}\right), \quad \text { where } W=\overline{\operatorname{span}}\left\{\psi(\cdot-k): k \in \mathbb{Z}^{n}, \psi \in \Psi\right\}
$$

It turns out that every semi-orthogonal tight framelet comes from a GMRA. Indeed, for a finite family $\Psi \subset L^{2}\left(\mathbb{R}^{n}\right)$ we define its space of negative dilates $V$ by

$$
V=\overline{\operatorname{span}}\left\{\psi_{j, k}: j<0, k \in \mathbb{Z}^{n}, \psi \in \Psi\right\} .
$$

We say that a framelet $\Psi$ is associated with a GMRA (or that it generates a GMRA) if its space of negative dilates $V$ satisfies (M1)-(M4). It is not hard to check that if $\Psi$ is a semi-orthogonal tight framelet then conditions (M1)-(M4) hold and, therefore, $V$ is a core space of a GMRA.

## 2. Scaling vectors for GMRA

The main goal of this section is to provide a constructive procedure for selecting a suitable set of generators $\Phi$ for the core space $V_{0}$ of a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$. The result of this procedure is a collection of (mutually orthogonal) quasi-orthogonal generators $\Phi=\left(\varphi_{j}\right)_{j \in J}$ called a scaling vector for $V_{0}$. (In particular, if $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is a usual MRA, we obtain a single orthogonal generator of $V_{0}$, usually called a scaling function.) Quasi-orthogonality means that integer shifts of the generator form a tight frame for the corresponding SI space. Such a space, that is generated by just one function, say $\varphi$, is called a principal shift-invariant (PSI) space and is denoted by $\mathcal{S}(\varphi)$. First, we shall review some basic results about SI spaces and the dimension and spectral functions.

Every shift-invariant space $V \subset L^{2}\left(\mathbb{R}^{n}\right)$ has a set of generators $\Phi$, that is, a countable family of functions whose integer shifts form a tight frame (with constant 1) for $V$, see [10, Theorem 3.3]. Although this family is not unique, the function

$$
\sigma_{V}(\xi)=\sum_{\varphi \in \Phi}|\hat{\varphi}(\xi)|^{2}
$$

does not depend (except on a set of measure zero) on the choice of the family of generators, see [15, Lemma 2.3]. Here, the Fourier transform is defined by

$$
\hat{f}(\xi)=\int_{\mathbb{R}^{n}} f(x) e^{-2 \pi i\langle x, \xi\rangle} d x
$$

We call $\sigma_{V}$ the spectral function of $V$. This notion was introduced by the authors in [15]. The basic property of $\sigma$ is that it is additive on countable orthogonal sums and that $\sigma_{L^{2}\left(\mathbb{R}^{n}\right)}=1$. The spectral function also behaves nicely under dilations since $\sigma_{D(V)}(\xi)=\sigma_{V}\left(B^{-1} \xi\right)$. Moreover, if $V$ is generated by a single function $\varphi$ then

$$
\sigma_{V}(\xi)= \begin{cases}|\hat{\varphi}(\xi)|^{2}\left(\sum_{k \in \mathbb{Z}^{n}}|\hat{\varphi}(\xi+k)|^{2}\right)^{-1} & \text { for } \xi \in \operatorname{supp} \hat{\varphi} \\ 0 & \text { otherwise }\end{cases}
$$

There are several other equivalent ways of defining the spectral function. The original one involves the range function $\mathcal{J}$, that is, a mapping from the torus $\mathbb{T}^{n}$ to the set of closed subspaces of $\ell^{2}\left(\mathbb{Z}^{n}\right)$. It turns that there is a $1-1$ correspondence between SI spaces and measurable range functions $\mathcal{J}$ given by

$$
V=\left\{f \in L^{2}\left(\mathbb{R}^{n}\right): \mathcal{T} f(\xi) \in \mathcal{J}(\xi) \text { for a.e. } \xi \in \mathbb{T}^{n}\right\}
$$

see [10, Proposition 1.5]. Here, $\mathcal{T}: L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{T}^{n}, \ell^{2}\left(\mathbb{Z}^{n}\right)\right)$ is an isometric isomorphism given by $\mathcal{T} f(\xi)=(\hat{f}(\xi+k))_{k \in \mathbb{Z}^{n}}$, where $\mathbb{T}^{n}$ is identified with $[-1 / 2,1 / 2)^{n}$. The spectral function $\sigma_{V}$ can be equivalently defined by

$$
\sigma_{V}(\xi+k)=\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2} \quad \text { for } \xi \in \mathbb{T}^{n} \text { and } k \in \mathbb{Z}^{n}
$$

where $\left\{e_{k}\right\}_{k \in \mathbb{Z}^{n}}$ denotes the standard basis of $\ell^{2}\left(\mathbb{Z}^{n}\right)$ and $P_{\mathcal{J}}(\xi)$ is an orthogonal projection of $\ell^{2}\left(\mathbb{Z}^{n}\right)$ onto $\mathcal{J}(\xi)$.

The spectral function also allows us to define the dimension function of $V$,

$$
\operatorname{dim}_{V}(\xi)=\sum_{k \in \mathbb{Z}^{n}} \sigma_{V}(\xi+k)
$$

The dimension function (also called the multiplicity function) is integer-valued and additive on countable orthogonal sums as well. Moreover, the minimal number of functions needed to generate $V$ is equal to the $L^{\infty}$ norm of $\operatorname{dim}_{V}$. Again, we refer the reader to $[10,15]$ for the proofs of all these facts.

The main feature of our generator selecting procedure is that it distinguishes the first generator $\varphi_{1}$ as having a dominating effect on all remaining generators. More precisely, the first generator $\varphi_{1}$ is chosen so that it exhausts the entire spectrum of the core space $V_{0}$ in some neighborhood of the origin. The fact that the space $V_{0}$ is refinable and this exhaustion property of $\varphi_{1}$ leads to a very special form of a matrix mask of $\Phi$, whose first column has zeros in every, but the first entry, near the origin, see Theorem 2.2.

Our procedure is somewhat reminiscent of the superfunction theory in the study of finitely generated shift-invariant (FSI) spaces by de Boor, DeVore, and Ron [22,23]. Among other things, the authors proved [23, Result 1.2] that an approximation order of an FSI space can be realized by some PSI space generated by a single function $\psi$, called a "superfunction." Therefore, $\psi$ has a dominating effect by providing the same approximation order as the whole FSI space generated by some finite collection of generators. This is analogous to our construction, where the first generator $\varphi_{1}$ makes all other generators to be innocuous near the origin in the Fourier domain, thus producing a special form of a matrix mask.

To achieve the above dominating effect we show the existence of a quasi-generator $\varphi_{0}$ of an SI space $V_{0}$ having the same spectral function as that of $V_{0}$ in a pre-specified localized region of the Fourier domain. Therefore, the generator $\varphi_{0}$ exhausts locally the space $V_{0}$ in that region.

Lemma 2.1. Assume that $V_{0} \subset L^{2}\left(\mathbb{R}^{n}\right)$ is SI. Let $K$ be a measurable subset of $\mathbb{R}^{n}$ such that

$$
|K \cap(K+l)|=0 \quad \text { for all } l \in \mathbb{Z}^{n} \backslash\{0\} .
$$

Let $\varphi=P_{V_{0}}\left(\check{\mathbf{1}}_{K}\right)$, where $P_{V_{0}}$ is the orthogonal projection onto $V_{0}$, and $\check{\mathbf{1}}_{K}$ is the inverse Fourier transform of the characteristic function $\mathbf{1}_{K}$. Define $\varphi_{0} \in L^{2}\left(\mathbb{R}^{n}\right)$ by

$$
\hat{\varphi}_{0}(\xi)= \begin{cases}\hat{\varphi}(\xi)\left(\sum_{k \in \mathbb{Z}^{n}}|\hat{\varphi}(\xi+k)|^{2}\right)^{-1 / 2}, & \xi \in \operatorname{supp} \hat{\varphi}  \tag{2.1}\\ 0, & \text { otherwise }\end{cases}
$$

Then $\varphi_{0} \in V_{0}$ is a quasi-orthogonal generator of $\mathcal{S}\left(\varphi_{0}\right)$ and

$$
\begin{equation*}
\sigma_{\mathcal{S}\left(\varphi_{0}\right)}(\xi)=\left|\hat{\varphi}_{0}(\xi)\right|^{2}=\sigma_{V_{0}}(\xi) \quad \text { for a.e. } \xi \in K \tag{2.2}
\end{equation*}
$$

Proof. Let $\varphi_{K} \in L^{2}\left(\mathbb{R}^{n}\right)$ be given by $\hat{\varphi}_{K}=\mathbf{1}_{K}$, and hence $\varphi=P_{V_{0}} \varphi_{K}$. Clearly, $\varphi_{0}$ is a quasiorthogonal generator of the PSI space $\mathcal{S}\left(\varphi_{0}\right)=\mathcal{S}(\varphi) \subset V_{0}$. In particular,

$$
\left|\hat{\varphi}_{0}(\xi)\right|^{2}=\sigma_{\mathcal{S}\left(\varphi_{0}\right)}(\xi)=\sigma_{\mathcal{S}(\varphi)}(\xi)
$$

Let $\mathcal{J}$ be the range function of $V_{0}$ with the corresponding orthogonal projections $P_{\mathcal{J}}(\xi)$. Then for any $f \in L^{2}\left(\mathbb{R}^{n}\right)$ we have

$$
\mathcal{T}\left(P_{V_{0}} f\right)(\xi)=P_{\mathcal{J}}(\xi)(\mathcal{T} f(\xi)) \quad \text { for a.e. } \xi \in \mathbb{T}^{n}
$$

Hence, for a.e. $\xi \in \mathbb{T}^{n}$,

$$
\mathcal{T} \varphi(\xi)=\mathcal{T}\left(P_{V_{0}} \varphi_{K}\right)(\xi)=P_{\mathcal{J}}(\xi)\left(\mathcal{T} \varphi_{K}(\xi)\right)= \begin{cases}P_{\mathcal{J}}(\xi) e_{k}, & \xi+k \in K, k \in \mathbb{Z}^{n} \\ 0, & \text { otherwise }\end{cases}
$$

Fix $k \in \mathbb{Z}^{n}$. If $\xi+k \in K, \xi \in \mathbb{T}^{n}$, and $\mathcal{T} \varphi(\xi) \neq 0$, then we necessarily have

$$
\begin{aligned}
\sigma_{\mathcal{S}(\varphi)}(\xi+k)=\frac{|\hat{\varphi}(\xi+k)|^{2}}{\|\mathcal{T} \varphi(\xi)\|^{2}} & =\frac{\left|\left\langle\mathcal{T} \varphi(\xi), e_{k}\right\rangle\right|^{2}}{\|\mathcal{T} \varphi(\xi)\|^{2}}=\frac{\left|\left\langle P_{\mathcal{J}}(\xi) e_{k}, e_{k}\right\rangle\right|^{2}}{\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2}}=\frac{\left|\left\langle P_{\mathcal{J}}(\xi)^{2} e_{k}, e_{k}\right\rangle\right|^{2}}{\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2}} \\
& =\frac{\left|\left\langle P_{\mathcal{J}}(\xi) e_{k}, P_{\mathcal{J}}(\xi) e_{k}\right\rangle\right|^{2}}{\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2}}=\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2}=\sigma_{V_{0}}(\xi+k)
\end{aligned}
$$

On the other hand, if $\xi+k \in K, \xi \in \mathbb{T}^{n}$, and $\mathcal{T} \varphi(\xi)=0$, then

$$
\sigma_{\mathcal{S}(\varphi)}(\xi+k)=0=\left\|P_{\mathcal{J}}(\xi) e_{k}\right\|^{2}=\sigma_{V_{0}}(\xi+k)
$$

Since $k \in \mathbb{Z}^{n}$ is arbitrary, this proves (2.2).

The next result provides a decomposition of any SI space $V_{0}$ as an orthogonal sum of carefully chosen PSI spaces, such that the first PSI space exhausts the spectrum of $V_{0}$ near the origin.

Theorem 2.1. Assume that $V_{0} \subset L^{2}\left(\mathbb{R}^{n}\right)$ is SI. Then there exists an orthogonal decomposition

$$
\begin{equation*}
V_{0}=\bigoplus_{j=1}^{\infty} \mathcal{S}\left(\varphi_{j}\right) \tag{2.3}
\end{equation*}
$$

such that each $\varphi_{j}$ is a quasi-orthogonal generator of $\mathcal{S}\left(\varphi_{j}\right)$ and

$$
\begin{equation*}
\text { supp } \operatorname{dim}_{\mathcal{S}\left(\varphi_{j}\right)}=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V_{0}}(\xi) \geqslant j\right\} \quad \text { for every } j \in \mathbb{N} . \tag{2.4}
\end{equation*}
$$

Furthermore, the spectral function of $\mathcal{S}\left(\varphi_{1}\right)$ coincides with that of $V_{0}$ near the origin, i.e.,

$$
\begin{equation*}
\sigma_{\mathcal{S}\left(\varphi_{1}\right)}(\xi)=\left|\hat{\varphi}_{1}(\xi)\right|^{2}=\sigma_{V_{0}}(\xi) \quad \text { for a.e. } \xi \in \mathbb{T}^{n} \tag{2.5}
\end{equation*}
$$

Proof. The existence of a decomposition satisfying (2.3) and (2.4) is already known, see [10]. The novelty of Theorem 2.1 lies in the fact that the first quasi-orthogonal generator $\varphi_{1}$ can be chosen to satisfy (2.5).

Let $\varphi_{0} \in V_{0}$ be a quasi-orthogonal generator guaranteed by Lemma 2.1 with $K=\mathbb{T}^{n}$. That is,

$$
\sigma_{\mathcal{S}\left(\varphi_{0}\right)}(\xi)=\left|\hat{\varphi}_{0}(\xi)\right|^{2}=\sigma_{V_{0}}(\xi) \quad \text { for a.e. } \xi \in \mathbb{T}^{n}
$$

Define $E=\operatorname{supp} \operatorname{dim}_{V_{0}} \backslash \operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\varphi_{0}\right)}$. Consider two possible cases. If $|E|>0$, then define an SI space $V=V_{0} \cap \check{L}^{2}(E)$. Here,

$$
\check{L}^{2}(E)=\left\{f \in L^{2}\left(\mathbb{R}^{n}\right): \operatorname{supp} \hat{f} \subset E\right\}
$$

Let $\varphi$ be a quasi-orthogonal generator of $V$ such that

$$
\text { supp } \operatorname{dim}_{\mathcal{S}(\varphi)}=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V}(\xi) \geqslant 1\right\}=\left\{\xi \in E: \operatorname{dim}_{V_{0}}(\xi) \geqslant 1\right\}=E
$$

Since $\varphi_{0} \in \check{L}^{2}\left(\mathbb{R}^{n} \backslash E\right), \varphi \in \check{L}^{2}(E)$, and the set $E$ is invariant under translations by $\mathbb{Z}^{n}, \varphi_{1}=$ $\varphi_{0}+\varphi$ is a quasi-orthogonal generator of $\mathcal{S}\left(\varphi_{1}\right)$. Moreover, $\varphi_{1} \in V_{0}$ since $\varphi_{0}, \varphi \in V_{0}$, and

$$
\text { supp } \operatorname{dim}_{\mathcal{S}\left(\varphi_{1}\right)}=\operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\varphi_{0}\right)} \cup \operatorname{supp}_{\operatorname{dim}}^{\mathcal{S}(\varphi)}\left(=\operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\varphi_{0}\right)} \cup E=\operatorname{supp} \operatorname{dim}_{V_{0}}\right.
$$

Hence, (2.4) holds for $j=1$. Since $\mathcal{S}\left(\varphi_{0}\right) \subset \mathcal{S}\left(\varphi_{1}\right) \subset V_{0}$, we have that for a.e. $\xi \in \mathbb{T}^{n}$,

$$
\sigma_{V_{0}}(\xi) \leqslant \sigma_{\mathcal{S}\left(\varphi_{0}\right)}(\xi) \leqslant \sigma_{\mathcal{S}\left(\varphi_{1}\right)}(\xi) \leqslant \sigma_{V_{0}}(\xi)
$$

which proves (2.5). In the case of $|E|=0$, we let $\varphi_{1}=\varphi_{0}$. Trivially, (2.4) holds for $j=1$ and (2.5) also holds.

Finally, it suffices to consider an SI space $V_{0} \ominus \mathcal{S}\left(\varphi_{1}\right)$ and its decomposition guaranteed by the first part of Theorem 2.1. That is we have

$$
\begin{gathered}
V_{0} \ominus \mathcal{S}\left(\varphi_{1}\right)=\bigoplus_{j=2}^{\infty} \mathcal{S}\left(\varphi_{j}\right) \\
\operatorname{supp}_{\operatorname{dim}}^{\mathcal{S}\left(\varphi_{j}\right)} \\
=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V_{0} \ominus \mathcal{S}\left(\varphi_{1}\right)}(\xi) \geqslant j-1\right\}=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V_{0}}(\xi) \geqslant j\right\} \quad \text { for } j \geqslant 2
\end{gathered}
$$

Therefore, $\left\{\varphi_{j}\right\}_{j=1}^{\infty}$ is the sequence of quasi-orthogonal generators fulfilling (2.3)-(2.5).
Theorem 2.1 leads naturally to the definition of an exhausting quasi-orthogonal vector for general SI spaces and an exhausting scaling vector for refinable SI spaces.

Definition 2.1. Suppose that $V_{0}$ is SI and for $j \in \mathbb{N}$ let

$$
\begin{equation*}
S_{j}=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V_{0}}(\xi) \geqslant j\right\} \tag{2.6}
\end{equation*}
$$

Let $J=\left\{j \in \mathbb{N}:\left|S_{j}\right|>0\right\}$. Naturally,

$$
J= \begin{cases}\{1, \ldots, L\} & \text { if } L=\operatorname{ess} \sup _{\xi \in \mathbb{R}^{n}} \operatorname{dim}_{V_{0}}(\xi)<\infty  \tag{2.7}\\ \mathbb{N} & \text { otherwise }\end{cases}
$$

A quasi-orthogonal vector for $V_{0}$ is defined as

$$
\Phi=\left(\varphi_{j}\right)_{j \in J},
$$

where $\left\{\varphi_{j}\right\}_{j \in J}$ are quasi-orthogonal generators as in Theorem 2.1 satisfying (2.3) and (2.4) only. In addition, if (2.5) holds, then we say that $\Phi$ is an exhausting quasi-orthogonal vector for $V_{0}$. The Fourier transform of $\Phi$,

$$
\hat{\Phi}(\xi)=\left(\hat{\varphi}_{j}(\xi)\right)_{j \in J}
$$

is treated as a column vector with values in $\ell^{2}(J)$, since

$$
\begin{equation*}
\|\hat{\Phi}(\xi)\|_{\ell^{2}(J)}^{2}=\sigma_{V_{0}}(\xi) \leqslant 1 . \tag{2.8}
\end{equation*}
$$

Also, define the diagonal matrix function of $V_{0}$ as

$$
\Omega(\xi)=\left[\begin{array}{cccc}
\mathbf{1}_{S_{1}}(\xi) & 0 & 0 & \cdots  \tag{2.9}\\
0 & \mathbf{1}_{S_{2}}(\xi) & 0 & \ldots \\
0 & 0 & \mathbf{1}_{S_{3}}(\xi) & \ldots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

Suppose that $\Phi=\left(\varphi_{j}\right)_{j \in J}$ is a quasi-orthogonal vector for $V_{0}$. Since each $\varphi_{j}$ is a quasiorthogonal generator of $\mathcal{S}\left(\varphi_{j}\right)$ and $\mathcal{S}\left(\varphi_{j}\right) \perp \mathcal{S}\left(\varphi_{j^{\prime}}\right)$ for $j \neq j^{\prime}$, we have that for a.e. $\xi \in \mathbb{R}^{n}$,

$$
\begin{aligned}
& \sum_{k \in \mathbb{Z}^{n}}\left|\hat{\varphi}_{j}(\xi+k)\right|^{2}=\mathbf{1}_{S_{j}}(\xi), \\
& \sum_{k \in \mathbb{Z}^{n}} \hat{\varphi}_{j}(\xi+k) \overline{\hat{\varphi}_{j^{\prime}}(\xi+k)}=0 \quad \text { for } j \neq j^{\prime} .
\end{aligned}
$$

Hence, in short

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}} \hat{\Phi}(\xi+k) \hat{\Phi}^{*}(\xi+k)=\Omega(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{2.10}
\end{equation*}
$$

Definition 2.2. Suppose that an SI space $V_{0}$ is refinable, that is $V_{0} \subset D\left(V_{0}\right)$. In this case a quasiorthogonal vector $\Phi=\left(\varphi_{j}\right)_{j \in J}$ for $V_{0}$ is called a scaling vector for $V_{0}$. In addition, if (2.5) holds, then $\Phi$ is an exhausting scaling vector for $V_{0}$.

The next result provides a characterization of elements of an SI space in terms of its quasiorthogonal vector. Proposition 2.1 is an immediate consequence of the corresponding wellknown result for PSI spaces; for example, see [27, Theorem 5.9].

Proposition 2.1. Suppose that an SI space $V_{0}$ is decomposed as in (2.3) and each $\varphi_{j}$ is a quasiorthogonal generator of $\mathcal{S}\left(\varphi_{j}\right)$. Then $f \in V_{0}$ if and only if

$$
\begin{equation*}
\hat{f}(\xi)=\sum_{j \in \mathbb{N}} r_{j}(\xi) \hat{\varphi}_{j}(\xi) \tag{2.11}
\end{equation*}
$$

where convergence is in $L^{2}$, each $r_{j}$ is $\mathbb{Z}^{n}$-periodic function in $L^{2}\left(S_{j}\right)$, and

$$
\|f\|^{2}=\sum_{j \in J}\left\|r_{j}\right\|^{2}
$$

Moreover, the sequence $\left\{r_{j}\right\}_{j \in \mathbb{N}}$ of such functions is unique.
Consequently, note that the series (2.11) converges a.e. after choosing a suitable subsequence. In particular, if the fibers of the SI space $V_{0}$ are finitely dimensional, meaning that (2.12) holds, the convergence in (2.11) is also in the almost everywhere sense. This observation leads to a simple characterization of refinability of such SI spaces.

Lemma 2.2. Suppose that $V_{0}$ is an SI space such that

$$
\begin{equation*}
\operatorname{dim}_{V_{0}}(\xi)<\infty \quad \text { for a.e. } \xi \tag{2.12}
\end{equation*}
$$

Suppose that $\Phi$ is a quasi-orthogonal vector of $V_{0}$, and $\left\{S_{j}\right\}_{j \in J}$ is given by (2.6) with J as in (2.7). Then the space $V_{0}$ is refinable with respect to the dilation $A$ if and only if

$$
\begin{equation*}
\hat{\Phi}(B \xi)=M(\xi) \hat{\Phi}(\xi) \tag{2.13}
\end{equation*}
$$

where $B=A^{T}$ and $M$ is $\mathbb{Z}^{n}$-periodic matrix function with entries $m_{i, j} \in L^{2}\left(S_{j}\right), i, j \in J$. Moreover, if such an $M$ exists, then it is unique.

Proof. Note that condition (2.12) guarantees that $\hat{\Phi}(\xi)$ has finitely many non-zero entries and the matrix product in (2.13) is meaningful. By Proposition 2.1, (2.13) implies that each $\varphi_{j} \in D\left(V_{0}\right)$. Since $D\left(V_{0}\right)$ is SI, we must have $V_{0} \subset D\left(V_{0}\right)$ and $V_{0}$ is refinable.

Conversely, if $V_{0}$ is refinable then the matrix $M$ satisfying (2.13) is uniquely determined with the use of Proposition 2.1 for $f=D^{-1} \varphi_{j}, j \in \mathbb{N}$.

A matrix function $M$ satisfying (2.13) is often called a matrix mask function of $\Phi$ or a low-pass matrix mask. We are now ready to prove the main result of this section providing a description of a matrix mask corresponding to an exhausting scaling vector $\Phi$ for $V_{0}$.

Theorem 2.2. Suppose that $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is a GMRA such that (2.12) holds. Let $\Phi$ be an exhausting scaling vector for $V_{0}$ and $M$ be the matrix mask function as in Lemma 2.2. Then

$$
\begin{equation*}
\sum_{d \in \mathcal{D}} M(\xi+d) M^{*}(\xi+d)=\Omega(B \xi) \tag{2.14}
\end{equation*}
$$

where $\mathcal{D}$ consists of representatives of distinct cosets of $B^{-1} \mathbb{Z}^{n} / \mathbb{Z}^{n}$. Moreover, the first column of $M(\xi)$ has zeros in every, but the first entry, near the origin in the sense that for a.e. $\xi \in \mathbb{R}^{n}$, there exists $N=N(\xi)$ such that

$$
\begin{equation*}
m_{i, 1}\left(B^{-j} \xi\right)=0 \quad \text { for } i \geqslant 2, j>N \tag{2.15}
\end{equation*}
$$

Furthermore, the upper-left corner of $M(\xi)$ has absolute value "almost equal" to 1 near the origin, i.e.,

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \prod_{j=k}^{\infty}\left|m_{1,1}\left(B^{-j} \xi\right)\right|=1 \quad \text { for a.e. } \xi \in \mathbb{R}^{n} . \tag{2.16}
\end{equation*}
$$

Proof. The condition on the support of $M$ implies that $M(\xi) \Omega(\xi)=M(\xi)$. Hence, by (2.10),

$$
\begin{aligned}
\Omega(B \xi) & =\sum_{k \in \mathbb{Z}^{n}} \hat{\Phi}(B \xi+k) \hat{\Phi}^{*}(B \xi+k) \\
& =\sum_{k \in \mathbb{Z}^{n}} M\left(\xi+B^{-1} k\right) \hat{\Phi}\left(\xi+B^{-1} k\right) \hat{\Phi}^{*}\left(\xi+B^{-1} k\right) M^{*}\left(\xi+B^{-1} k\right) \\
& =\sum_{d \in \mathcal{D}} M(\xi+d) \Omega(\xi+d) M^{*}(\xi+d)=\sum_{d \in \mathcal{D}} M(\xi+d) M^{*}(\xi+d),
\end{aligned}
$$

which proves (2.14).
To show (2.15) we will use the fact [15, Lemma 2.7] that

$$
\begin{equation*}
\lim _{j \rightarrow \infty}\left|\hat{\varphi}_{1}\left(B^{-j} \xi\right)\right|^{2}=\lim _{j \rightarrow \infty} \sigma_{V_{0}}\left(B^{-j} \xi\right)=\lim _{j \rightarrow \infty} \sigma_{V_{j}}(\xi)=1 \quad \text { a.e. } \xi \in \mathbb{R}^{n} \tag{2.17}
\end{equation*}
$$

Combining this with (2.8) yields that there exists $N=N(\xi)$ such that the first coordinate of $\hat{\Phi}\left(B^{-j} \xi\right)$ is non-zero and all others are zero for all $j \geqslant N$. By (2.13),

$$
\hat{\varphi}_{i}\left(B^{-j+1} \xi\right)=m_{i, 1}\left(B^{-j} \xi\right) \hat{\varphi}_{1}\left(B^{-j} \xi\right) \quad \text { for } j \geqslant N, i \in \mathbb{N} .
$$

Since $\hat{\varphi}_{i}\left(B^{-j+1} \xi\right)=0$ for $i \geqslant 2$ and $\hat{\varphi}_{1}\left(B^{-j} \xi\right) \neq 0$ for $j>N$, we have (2.15).
To show (2.16) it suffices to observe that for every $l>k \geqslant N$,

$$
\hat{\varphi}_{1}\left(B^{-k} \xi\right)=\hat{\varphi}_{1}\left(B^{-l} \xi\right) \prod_{j=k+1}^{l} m_{1,1}\left(B^{-j} \xi\right) .
$$

By (2.17),

$$
\left|\hat{\varphi}_{1}\left(B^{-k} \xi\right)\right|=\lim _{l \rightarrow \infty}\left|\hat{\varphi}_{1}\left(B^{-l} \xi\right)\right| \prod_{j=k+1}^{l}\left|m_{1,1}\left(B^{-j} \xi\right)\right|=\prod_{j=k+1}^{\infty}\left|m_{1,1}\left(B^{-j} \xi\right)\right|,
$$

which proves (2.16) by letting $k \rightarrow \infty$.
The fact that $M$ satisfies the condition (2.14), which is also called a generalized low-pass filter or generalized conjugate mirror filter, is due Baggett, Courter, and Merrill [3, Theorem 2.3]. This condition holds for all scaling vectors $\Phi$ (not necessarily exhausting) and it is an analogue of the usual quadrature-mirror equation (1.1). The additional assumption that $\Phi$ is an exhausting scaling vector implies that the first column of $M$ must be of a special form ( $m_{1,1}(\xi), 0,0, \ldots$ ) with $\left|m_{1,1}(\xi)\right| \approx 1$ for $\xi$ near 0 . Moreover, (2.14) implies that

$$
\sum_{d \in \mathcal{D}}\left|m_{1,1}(\xi+d)\right|^{2} \leqslant \mathbf{1}_{S_{1}}(B \xi) \leqslant 1 \quad \text { for a.e. } \xi \in \mathbb{T}^{n}
$$

For these reasons $m_{1,1}(\xi)$ plays a role similar to that of a usual low-pass filter and it has a dominating effect on the entire matrix mask function $M$. These issues are further explored in Section 4, where the procedure for reconstructing a scaling vector from its matrix mask is presented. We should also emphasize that conditions (2.14)-(2.15) are only necessary and not sufficient for guaranteeing that $M$ is a matrix mask of some scaling vector. This is a simple consequence of the usual MRA case, where (1.1) and (2.16) alone are not enough to produce the scaling function and some extra conditions, such as Lawton's or Cohen's conditions are needed [25,27].

Next, we will look at semi-orthogonal wavelets associated to a GMRA. In [16] we pointed out that one can always find a semi-orthogonal wavelet (possibly with infinite number of generators) associated to any GMRA. To be more precise, let us state the following

Theorem 2.3. Suppose that $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is a GMRA such that (2.12) holds. Then there exists a semiorthogonal wavelet $\left(\psi^{j}\right)_{j \in \tilde{J}} \subset L^{2}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{gather*}
W_{0}:=V_{1} \ominus V_{0}=\bigoplus_{j \in \tilde{J}} \mathcal{S}\left(\psi_{j}\right)  \tag{2.18}\\
\operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\psi_{j}\right)}=\tilde{S}_{j}:=\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{W_{0}}(\xi) \geqslant j\right\} .
\end{gather*}
$$

Here, $\tilde{J}$ is either $\{1, \ldots, N\}$ or $\mathbb{N}$.

Conversely, suppose that we have a semi-orthogonal wavelet $\Psi=\left(\psi^{\tilde{j}}\right)_{\tilde{j} \in \tilde{J}}$, where $\tilde{J}=$ $\{1, \ldots, N\}$ is finite, which is associated with a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$; that is (2.18) holds. Equivalently, a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ associated to $\Psi$ is given by

$$
V_{j}=\operatorname{span}\left\{D^{i} T_{k} \psi^{\tilde{j}}: i<j, k \in \mathbb{Z}^{n}, \tilde{j} \in \tilde{J}\right\} \quad \text { for } j \in \mathbb{Z}
$$

Let $\Psi$ be the column vector defined as $\Psi=\left(\psi^{j}\right)_{j \in \tilde{J}}$. By Proposition 2.1, there exists a matrix function $H(\xi)=\left(h_{i, j}(\xi)\right)_{i \in \tilde{J}, j \in J}$ such that

$$
\hat{\Psi}(B \xi)=H(\xi) \hat{\Phi}(\xi)
$$

and $h_{i, j} \in L^{2}\left(S_{j}\right)$. Let $\tilde{\Omega}$ be the diagonal matrix function corresponding to $\Psi$, i.e.,

$$
\begin{equation*}
\tilde{\Omega}(\xi)=\operatorname{diag}\left\{\mathbf{1}_{\tilde{S}_{j}}(\xi): j \in \tilde{J}\right\}, \quad \text { where } \tilde{S}_{j}=\operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\psi^{j}\right)} \tag{2.19}
\end{equation*}
$$

Then we have the following description of a matrix mask function $H$ corresponding to a semiorthogonal wavelet $\Psi$, called a high-pass matrix mask or complementary conjugate mirror filter in [3].

Proposition 2.2. Suppose $\Psi$ is a semi-orthogonal wavelet associated with a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$. Let $M$ and $H$ be the low-pass and high-pass matrix mask functions as above. Then

$$
\begin{gather*}
\sum_{d \in \mathcal{D}} H(\xi+d) H^{*}(\xi+d)=\tilde{\Omega}(B \xi),  \tag{2.20}\\
\sum_{d \in \mathcal{D}} M(\xi+d) H^{*}(\xi+d)=\sum_{d \in \mathcal{D}} H(\xi+d) M^{*}(\xi+d)=0 . \tag{2.21}
\end{gather*}
$$

Proof. The condition on the support of $H$ implies that $H(\xi) \Omega(\xi)=H(\xi)$. Hence, by (2.10),

$$
\begin{aligned}
\tilde{\Omega}(B \xi) & =\sum_{k \in \mathbb{Z}^{n}} \hat{\Psi}(B \xi+k) \hat{\Psi}^{*}(B \xi+k) \\
& =\sum_{k \in \mathbb{Z}^{n}} H\left(\xi+B^{-1} k\right) \hat{\Phi}\left(\xi+B^{-1} k\right) \hat{\Phi}^{*}\left(\xi+B^{-1} k\right) H^{*}\left(\xi+B^{-1} k\right) \\
& =\sum_{d \in \mathcal{D}} H(\xi+d) \Omega(\xi+d) H^{*}(\xi+d)=\sum_{d \in \mathcal{D}} H(\xi+d) H^{*}(\xi+d),
\end{aligned}
$$

which proves (2.20). Likewise,

$$
\begin{aligned}
0 & =\sum_{k \in \mathbb{Z}^{n}} \hat{\Phi}(B \xi+k) \hat{\Psi}^{*}(B \xi+k) \\
& =\sum_{k \in \mathbb{Z}^{n}} M\left(\xi+B^{-1} k\right) \hat{\Phi}\left(\xi+B^{-1} k\right) \hat{\Phi}^{*}\left(\xi+B^{-1} k\right) H^{*}\left(\xi+B^{-1} k\right) \\
& =\sum_{d \in \mathcal{D}} M(\xi+d) \Omega(\xi+d) H^{*}(\xi+d)=\sum_{d \in \mathcal{D}} M(\xi+d) H^{*}(\xi+d),
\end{aligned}
$$

which proves (2.21).

We can now summarize the procedure of generating low-pass $M$ and high-pass $H$ matrix masks as follows. Given a semi-orthogonal wavelet $\Psi$, we consider its associated GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$. Theorem 2.1 provides an exhausting scaling vector $\Phi$ for $V_{0}$. Then, Theorem 2.2 and Proposition 2.2 yield a low-pass matrix mask $M$ and a high-pass matrix mask $H$ satisfying (2.14)-(2.16) and (2.20)-(2.21), respectively. Since the decomposition of Theorem 2.1 is not unique, the above procedure yields a multitude of low-pass and high-pass matrix masks for a fixed semi-orthogonal wavelet. This fact makes the problem of reconstructing scaling vector and semi-orthogonal wavelet from their corresponding low-pass and high-pass matrix masks a highly non-trivial task. Fortunately, the exhausting property of $\Phi$ will counterbalance such inherent non-uniqueness. This will be explored in Sections 4 and 5, where the appropriate reconstruction procedure is provided, see Theorem 5.4.

We end this section with two remarks involving orthogonality of columns versus rows for combined low-pass and high-pass matrix masks. Remark 2.1 is a slight refinement of [4, Theorem 2.5], where matrix mask functions $M$ and $H$ were assumed to have finite size.

Remark 2.1. Let $T$ be the combined matrix mask function of a scaling vector $\Phi$ and the associated semi-orthogonal wavelet $\Psi$ given by

$$
T(\xi)=\left[\begin{array}{ccc}
M\left(\xi+d_{1}\right) & \ldots & M\left(\xi+d_{q}\right) \\
H\left(\xi+d_{1}\right) & \ldots & H\left(\xi+d_{q}\right)
\end{array}\right]
$$

where $d_{1}, \ldots, d_{q}$ are representatives of distinct cosets of $B^{-1} \mathbb{Z}^{n} / \mathbb{Z}^{n}$. Note that for a.e. $\xi \in \mathbb{T}^{n}, T(\xi)$ has only a finite number of non-zero entries. More precisely, there are only $\sum_{d \in \mathcal{D}} \operatorname{dim}_{V_{0}}(\xi+d)$ non-zero columns and $\operatorname{dim}_{V_{0}}(B \xi)+\operatorname{dim}_{W_{0}}(B \xi)$ non-zero rows. Let $\tilde{T}(\xi)$ be a finite sub-matrix of $T(\xi)$ consisting of only these columns and rows. The consistency equation of Baggett, see [6,7,15], says that

$$
\begin{equation*}
\operatorname{dim}_{V_{1}}(B \xi)=\sum_{d \in \mathcal{D}} \operatorname{dim}_{V_{0}}(\xi+d)=\operatorname{dim}_{V_{0}}(B \xi)+\operatorname{dim}_{W_{0}}(B \xi) \quad \text { for a.e. } \xi \in \mathbb{T}^{n} \tag{2.22}
\end{equation*}
$$

which implies that the matrix $\tilde{T}(\xi)$ is square for a.e. $\xi \in \mathbb{T}^{n}$. Moreover, Theorem 2.2 and Proposition 2.2 imply that the rows of the matrix $\tilde{T}(\xi)$ are mutually orthogonal and normalized, that is $\tilde{T}(\xi)$ is a unitary matrix. Consequently, the columns of $\tilde{T}(\xi)$ are mutually orthogonal and normalized.

Remark 2.2. One could consider the combined matrix mask function $T(\xi)$ corresponding to a more general situation when $\Psi$ is a framelet obtained by a similar procedure. In this case, the rows of $T(\xi)$ do not have to be mutually orthogonal, anymore. In fact, the sub-matrix of non-zero rows and columns $\tilde{T}(\xi)$ does not have to be square, since we can have many more generators in $\Psi$, and hence more rows in the matrix mask function $H(\xi)$. It turns out that unlike the situation of semi-orthogonal wavelets, where orthogonality of rows of $T(\xi)$ is necessary, orthogonality of columns plays a critical role for general framelets. This will be explored in the next section.

## 3. Unitary Extension Principle

The Unitary Extension Principle, and its generalizations such as Oblique Extension Principle, are powerful tools in constructing tight framelets [8,9,21,29]. Since these techniques are used
for constructing framelets with many desired properties such as smoothness, compact support, vanishing moments, etc., the Unitary Extension Principle is very often stated with some very mild and convenient regularity assumptions on a refinable function $\varphi$.

Since the interest of our work lies mainly in $L^{2}$ theory of framelets and wavelets, it is imperative to avoid any regularity assumptions, regardless of their mildness, limiting the applicability of our results. Furthermore, we are also forced to study situations where we are given a refinable vector consisting of infinite number of functions. Since these two problems were not adequately addressed yet, we provide an extension of Unitary Extension Principle, that is perfectly adapted to the $L^{2}$ theory. We start with a definition of a refinable vector consisting of potentially infinitely many functions.

Definition 3.1. We say that $\Phi=\left(\varphi_{j}\right)_{j \in J} \subset L^{2}\left(\mathbb{R}^{n}\right)$ is a refinable vector, where $J=\{1, \ldots, N\}$ or $J=\mathbb{N}$, if

$$
\begin{equation*}
\hat{\Phi}(B \xi)=M(\xi) \hat{\Phi}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{3.1}
\end{equation*}
$$

where $B=A^{T}$ and $M=\left(m_{i, j}\right)_{i, j \in J}$ is a matrix of $\mathbb{Z}^{n}$-periodic, measurable functions.
In order to make sense of (3.1) in the case when $J=\mathbb{N}$, we assume additionally that

$$
\begin{equation*}
\sum_{j \in J} \mathbf{1}_{R_{j}}(\xi)<\infty \quad \text { for a.e. } \xi \in \mathbb{T}^{n} \tag{3.2}
\end{equation*}
$$

where

$$
R_{j}=\operatorname{supp}_{\operatorname{dim}_{\mathcal{S}}\left(\varphi_{j}\right)} .
$$

Note that we can always assume that $\operatorname{supp} m_{i, j} \subset R_{j}$, since the values of $m_{i, j}$ outside of $R_{j}$ do not affect (3.1).

Remark 3.1. Condition (3.2) is a technical matter that allows us to talk meaningfully about Eq. (3.1). However, if the matrix $M(\xi)$ has only finitely many non-zero entries for a.e. $\xi \in \mathbb{R}^{n}$, then (3.1) makes sense right away. Moreover, in this simple case, (3.2) follows from (3.1).

Theorem 3.1 is a generalization of the Unitary Extension Principle of Ron and Shen [29] to a situation when a refinable vector $\Phi$ is infinite. We note that the original result of Ron and Shen, in the case when $\Phi$ is finite, requires certain mild decay assumptions on $\Phi$, see [21,29]. Nevertheless, Theorem 3.1 shows that these decay assumptions are unnecessary and they can be safely removed.

Theorem 3.1. Suppose $\Phi=\left(\varphi_{j}\right)_{j \in J}$ is a refinable vector with a mask $M$ such that

$$
\begin{equation*}
\sum_{j \in J}\left\|\varphi_{j}\right\|^{2}=\int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|_{\ell^{2}(J)}^{2} d \xi<\infty \tag{3.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{j \rightarrow \infty}\left\|\hat{\Phi}\left(B^{-j} \xi\right)\right\|=1 \quad \text { for a.e. } \xi \in \mathbb{R}^{n} . \tag{3.4}
\end{equation*}
$$

Suppose also that $\Psi=\left(\psi^{j}\right)_{j \in \tilde{J}}$, where $\tilde{J}=\{1, \ldots, N\}$ is finite, is given by

$$
\begin{equation*}
\hat{\Psi}(B \xi)=H(\xi) \hat{\Phi}(\xi) \tag{3.5}
\end{equation*}
$$

where $H=\left(h_{i, j}\right)_{i \in \tilde{J}, j \in J}$ is $\mathbb{Z}^{n}$-periodic, measurable matrix function satisfying

$$
\begin{equation*}
M^{*}(\xi) M(\xi+d)+H^{*}(\xi) H(\xi+d)=\Omega(\xi) \delta_{0, d} \quad \text { for a.e. } \xi \tag{3.6}
\end{equation*}
$$

and for any $d \in \mathcal{D}$.
Then $\Psi \subset L^{2}\left(\mathbb{R}^{n}\right)$ is a tight framelet.
Proof. It suffices to verify that $\Psi \subset L^{2}\left(\mathbb{R}^{n}\right)$ satisfies the characterization equations for tight framelets [11,24]

$$
\begin{gather*}
\sum_{j \in \mathbb{Z}}\left\|\hat{\Psi}\left(B^{j} \xi\right)\right\|^{2}=1 \quad \text { for a.e. } \xi  \tag{3.7}\\
\sum_{j=0}^{\infty} \hat{\Psi}^{*}\left(B^{j} \xi\right) \hat{\Psi}\left(B^{j}(\xi+q)\right)=0 \quad \text { for a.e. } \xi, \text { and all } q \in \mathbb{Z}^{n} \backslash B \mathbb{Z}^{n} . \tag{3.8}
\end{gather*}
$$

Note that for any $j \in \mathbb{Z}$,

$$
\begin{align*}
\left\|\hat{\Psi}\left(B^{j} \xi\right)\right\|^{2}+\left\|\hat{\Phi}\left(B^{j} \xi\right)\right\|^{2}= & \hat{\Psi}^{*}\left(B^{j} \xi\right) \hat{\Psi}\left(B^{j} \xi\right)+\hat{\Phi}^{*}\left(B^{j} \xi\right) \hat{\Phi}\left(B^{j} \xi\right) \\
= & \hat{\Phi}^{*}\left(B^{j-1} \xi\right) H^{*}\left(B^{j-1} \xi\right) H\left(B^{j-1} \xi\right) \hat{\Phi}\left(B^{j-1} \xi\right) \\
& +\hat{\Phi}^{*}\left(B^{j-1} \xi\right) M^{*}\left(B^{j-1} \xi\right) M\left(B^{j-1} \xi\right) \hat{\Phi}\left(B^{j-1} \xi\right) \\
= & \hat{\Phi}^{*}\left(B^{j-1} \xi\right) \Omega\left(B^{j-1} \xi\right) \hat{\Phi}\left(B^{j-1} \xi\right)=\left\|\hat{\Phi}\left(B^{j-1} \xi\right)\right\|^{2} \tag{3.9}
\end{align*}
$$

where in the last step we used that supp $\hat{\varphi}_{i} \subset S_{i}$. Therefore,

$$
\int_{\mathbb{R}^{n}}\|\hat{\Psi}(\xi)\|^{2} d \xi=(|\operatorname{det} A|-1) \int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2} d \xi<\infty
$$

and the fact that $\Psi \subset L^{2}\left(\mathbb{R}^{n}\right)$ is forced by (3.5) and (3.6).
Next, we claim that

$$
\begin{equation*}
\lim _{j \rightarrow \infty}\left\|\hat{\Phi}\left(B^{j} \xi\right)\right\|=0 \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{3.10}
\end{equation*}
$$

Otherwise, due to monotonicity of the sequence $\left(\left\|\hat{\Phi}\left(B^{j} \xi\right)\right\|\right)_{j \in \mathbb{Z}}$, we could find $\delta>0$ such that

$$
E=\left\{\xi \in \mathbb{R}^{n}:\left\|\hat{\Phi}\left(B^{j} \xi\right)\right\|>\delta \text { for all } j \in \mathbb{Z}\right\}
$$

has a positive measure. Since $B E=E, E$ must have infinite Lebesgue measure and consequently

$$
\int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2} d \xi \geqslant \int_{E}\|\hat{\Phi}(\xi)\|^{2} d \xi=\infty
$$

which contradicts (3.3). Thus, (3.10) holds and together with (3.4), (3.9) it implies (3.7).
Likewise for any $j \geqslant 1$ and $q \in \mathbb{Z}^{n} \backslash B \mathbb{Z}^{n}$,

$$
\begin{aligned}
\hat{\Psi}^{*} & \left(B^{j} \xi\right) \hat{\Psi}\left(B^{j}(\xi+q)\right)+\hat{\Phi}^{*}\left(B^{j} \xi\right) \hat{\Phi}\left(B^{j}(\xi+q)\right) \\
= & \hat{\Phi}^{*}\left(B^{j-1} \xi\right) H^{*}\left(B^{j-1} \xi\right) H\left(B^{j-1}(\xi+q)\right) \hat{\Phi}\left(B^{j-1}(\xi+q)\right) \\
& \quad+\hat{\Phi}^{*}\left(B^{j-1} \xi\right) M^{*}\left(B^{j-1} \xi\right) M\left(B^{j-1}(\xi+q)\right) \hat{\Phi}\left(B^{j-1}(\xi+q)\right) \\
= & \hat{\Phi}^{*}\left(B^{j-1} \xi\right) \Omega\left(B^{j-1} \xi\right) \hat{\Phi}\left(B^{j-1}(\xi+q)\right) \\
= & \hat{\Phi}^{*}\left(B^{j-1} \xi\right) \hat{\Phi}\left(B^{j-1}(\xi+q)\right),
\end{aligned}
$$

where in the penultimate step we used $\mathbb{Z}^{n}$-periodicity of $M$ and $H$. The same calculation for $j=0$ together with the observation that

$$
H^{*}\left(B^{-1} \xi\right) H\left(B^{-1}(\xi+q)\right)+M^{*}\left(B^{-1} \xi\right) H\left(B^{-1}(\xi+q)\right)=0
$$

yields that

$$
\hat{\Psi}^{*}(\xi) \hat{\Psi}(\xi+q)+\hat{\Phi}^{*}(\xi) \hat{\Phi}(\xi+q)=0 .
$$

Combining these identities with

$$
\lim _{j \rightarrow \infty}\left|\hat{\Phi}^{*}\left(B^{j} \xi\right) \hat{\Phi}\left(B^{j}(\xi+q)\right)\right| \leqslant \lim _{j \rightarrow \infty}\left\|\hat{\Phi}\left(B^{j} \xi\right)\right\|\left\|\hat{\Phi}\left(B^{j}(\xi+q)\right)\right\|=0 \quad \text { for a.e. } \xi \in \mathbb{R}^{n}
$$

proves (3.8).
Remark 3.2. The Unitary Extension Principle in the form of Theorem 3.1 yields not only a tight framelet $\Psi$ but also two GMRAs. Indeed, every function $\Phi$ satisfying (3.1)-(3.4) generates a GMRA with a core space $V_{0}$ generated by the integer shifts of the functions $\varphi_{j}, j \in J$. The other GMRA is the one, whose core space $\tilde{V}_{0}$ is the space of negative dilates of the tight framelet $\Psi$, see Theorem 6.1. While we always have $\tilde{V}_{0} \subset V_{0}$, it is not clear whether we have the converse inclusion, i.e., whether these two GMRAs are the same. This open question was raised by Baggett, Jorgensen, Merrill, and Packer in [4].

In the next section we will use Theorem 3.1 to give a general construction procedure of tight framelets. To this end, it is convenient to prove the following fact about functions satisfying an inequality reminiscent of Baggett's consistency equation.

Lemma 3.1. Suppose that $m: \mathbb{R}^{n} \rightarrow[0, \infty)$ is $\mathbb{Z}^{n}$-periodic, measurable function such that

$$
\begin{equation*}
\sum_{d \in D} m(\xi+d) \leqslant m(B \xi)+M \quad \text { for a.e. } \xi \in \mathbb{T}^{n} \tag{3.11}
\end{equation*}
$$

for some $M \geqslant 0$. Then $m$ is integrable over its period and

$$
\begin{equation*}
\int_{\mathbb{T}^{n}} m(\xi) d \xi \leqslant M /(|\operatorname{det} A|-1) \tag{3.12}
\end{equation*}
$$

Heuristically, Lemma 3.1 seems to be trivial. Integrating (3.11) over $\mathbb{T}^{n}$ yields

$$
|\operatorname{det} A| \int_{\mathbb{T}^{n}} m(\xi) d \xi \leqslant \int_{\mathbb{T}^{n}} m(\xi) d \xi+M
$$

Unfortunately, we do not know a priori whether $m$ is integrable and a much more complicated argument is necessary. Despite its simplicity, we could not find Lemma 3.1 in the existing literature and therefore we provide its proof.

Proof. For an integer $N \geqslant 0$, let $R_{N}(\xi)$ be the "Riemann sum" of $m$ of depth $N$ given by

$$
R_{N}(\xi)=R_{N}^{m}(\xi):=\frac{1}{|\operatorname{det} A|^{N}} \sum_{\epsilon_{0}, \ldots, \epsilon_{N-1} \in \mathcal{D}} m\left(\xi+\sum_{i=0}^{N-1} B^{-i} \epsilon_{i}\right)
$$

It is clear that $R_{N}(\xi)$ is measurable and $B^{-N} \mathbb{Z}^{n}$-periodic, since all the sums of the form $\sum_{i=0}^{N-1} B^{-i} \epsilon_{i}$, where $\epsilon_{0}, \ldots, \epsilon_{N-1} \in \mathcal{D}$, are representatives of distinct cosets of $B^{-N} \mathbb{Z}^{n} / \mathbb{Z}^{n}$. Here, $\mathcal{D}$ consists as usual of representatives of distinct cosets of $B^{-1} \mathbb{Z}^{n} / \mathbb{Z}^{n}$. By (3.11),

$$
|\operatorname{det} A| R_{N}(\xi) \leqslant R_{N-1}(B \xi)+M \quad \text { for any } N \geqslant 1 .
$$

Hence, by iteration,

$$
m\left(B^{N} \xi\right)=R_{0}\left(B^{N} \xi\right) \geqslant|\operatorname{det} A|^{N} R_{N}(\xi)-M \frac{|\operatorname{det} A|^{N}-1}{|\operatorname{det} A|-1}
$$

Take any $C>M /(|\operatorname{det} A|-1)$ and let $\delta=C-M /(|\operatorname{det} A|-1)$. Then

$$
\left\{\xi \in \mathbb{R}^{n}: R_{N}(\xi) \geqslant C\right\} \subset\left\{\xi \in \mathbb{R}^{n}: m\left(B^{N} \xi\right) \geqslant \delta|\operatorname{det} A|^{N}\right\} .
$$

For a fixed $K>0$, let $R_{N}^{\prime}(\xi)=R_{N}^{m^{\prime}}(\xi)$, where $m^{\prime}$ is a truncation of $m$ at height $K$ given by $m^{\prime}(\xi)=\min (m(\xi), K)$. It is clear that each $R_{N}^{\prime}(\xi)$ is $B^{-N} \mathbb{Z}^{n}$-periodic, measurable and bounded by $K$. Furthermore,

$$
\begin{align*}
\left|\left\{\xi \in \mathbb{T}^{n}: R_{N}^{\prime}(\xi) \geqslant C\right\}\right| & \leqslant\left|\left\{\xi \in \mathbb{T}^{n}: m\left(B^{N} \xi\right) \geqslant \delta|\operatorname{det} A|^{N}\right\}\right| \\
& =\left|\left\{\xi \in \mathbb{T}^{n}: m(\xi) \geqslant \delta|\operatorname{det} A|^{N}\right\}\right| \rightarrow 0 \quad \text { as } N \rightarrow \infty \tag{3.13}
\end{align*}
$$

Since $R_{N}^{\prime}$ 's are bounded, there exists a subsequence $\left\{N_{i}\right\}$ such that $\left\{R_{N_{i}}^{\prime}(\xi)\right\}$ converges pointwise a.e. to some $f(\xi)$. Since $f(\xi)$ must be periodic with respect to every lattice $B^{-N} \mathbb{Z}^{n} \subset B^{-N+1} \mathbb{Z}^{n}$,
and $\bigcup_{N=0}^{\infty} B^{-N} \mathbb{Z}^{n}$ is dense, $f(\xi)$ must be a constant function, $f(\xi)=C_{0}$. By (3.13), $0 \leqslant C_{0} \leqslant$ $M /(|\operatorname{det} A|-1)$. Moreover,

$$
C_{0}=\int_{\mathbb{T}^{n}} f(\xi) d \xi=\lim _{i \rightarrow \infty} \int_{\mathbb{T}^{n}} R_{N_{i}}^{\prime}(\xi) d \xi=\int_{\mathbb{T}^{n}} m^{\prime}(\xi) d \xi=\int_{\mathbb{T}^{n}} \min (m(\xi), K) d \xi
$$

Hence, letting $K \rightarrow \infty$ allows to obtain (3.12) by the monotone convergence theorem.

## 4. Construction of tight framelets

The main goal of this section is to provide a general reconstruction procedure for scaling vectors and semi-orthogonal wavelets from their corresponding low-pass and high-pass matrix masks. Hence, the goal is to reverse the flow of Section 2 by starting with a low-pass matrix mask function $M$ satisfying conditions (2.14)-(2.16). Theorem 2.2 shows that this is a perfectly reasonable assumption, since any matrix mask function of an exhausting scaling vector must satisfy them. The key ingredient of our approach is a rather complicated procedure yielding a refinable vector $\Phi$ corresponding to the mask $M$, see Theorem 4.2. This, combined with the Unitary Extension Principle and appropriate conditions on a high-pass mask $H$, yields a tight framelet $\Psi$ (see Theorem 4.3).

In general, we can only expect that $\Psi$ is a tight framelet. However, if we know a priori that our low-pass and high-pass matrix masks correspond to some semi-orthogonal wavelet $\Psi$, then we prove that our procedure is flexible enough to recover $\Psi$ itself. In particular, every orthogonal wavelet $\Psi$ can be obtained by our recovery procedure via low-pass and high-pass matrix masks manipulations. This will be shown in the following section.

To start the construction of tight framelets we must recall a characterization of the dimension function associated to a GMRA proved in [15].

Theorem 4.1. Suppose $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is a GMRA. Then the dimension function of the core space $V_{0}$, $m(\xi)=\operatorname{dim}_{V_{0}}(\xi)$, satisfies the following conditions:
(D1) $m: \mathbb{R}^{n} \rightarrow \mathbb{N} \cup\{0, \infty\}$ is a measurable $\mathbb{Z}^{n}$-periodic function;
(D2) $\sum_{d \in \mathcal{D}} m(\xi+d) \geqslant m(B \xi)$ for a.e. $\xi \in \mathbb{R}^{n}$;
(D3) $\sum_{k \in \mathbb{Z}^{n}} \mathbf{1}_{\Delta}(\xi+k) \geqslant m(\xi)$ for a.e. $\xi \in \mathbb{R}^{n}$, where

$$
\Delta=\left\{\xi \in \mathbb{R}^{n}: m\left(B^{-j} \xi\right) \geqslant 1 \text { for } j \in \mathbb{N} \cup\{0\}\right\} ;
$$

(D4) $\liminf _{j \rightarrow \infty} m\left(B^{-j} \xi\right) \geqslant 1$ for a.e. $\xi \in \mathbb{R}^{n}$.
Conversely, if $m$ satisfies (D1)-(D4), then there exists a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ such that $\operatorname{dim}_{V_{0}}(\xi)=m(\xi)$.

Our construction is based on a function $m$ that satisfies conditions (D1)-(D4) of the above theorem. However, to ensure the existence of a tight framelet we shall add two more assumptions. Namely,
(D5) $m \in L^{1}\left(\mathbb{T}^{n}\right)$;
and
(D6) $m$ is finite a.e. and there is $N \in \mathbb{N}$ such that for a.e. $\xi \in \mathbb{T}^{n}$ we have

$$
\sum_{d \in \mathcal{D}} m(\xi+d) \leqslant m(B \xi)+N
$$

To motivate these final conditions we include the following
Proposition 4.1. Let $\Psi=\left\{\psi^{1}, \ldots, \psi^{N}\right\} \subset L^{2}\left(\mathbb{R}^{n}\right)$ be a tight framelet and $V_{0}$ its space of negative dilates. If $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ forms a GMRA, then $m=\operatorname{dim}_{V_{0}}$ satisfies (D5) and (D6).

Proof. First we conduct the standard orthogonalization procedure. That is, for $j \in \mathbb{Z}$ we define $W_{j}:=V_{j+1} \ominus V_{j}$ and observe that

$$
\begin{equation*}
\bigoplus_{i \in \mathbb{T}} W_{j}=L^{2}\left(\mathbb{R}^{n}\right) \tag{4.1}
\end{equation*}
$$

Clearly, $W_{0}$ is a shift-invariant space generated by $\left\{\psi-P_{V_{0}} \psi\right\}_{\psi \in \Psi}$, where $P_{V_{0}}$ is the orthogonal projection on $V_{0}$. By Theorem 2.1 we can find quasi-orthogonal generators $\Phi=\left\{\varphi_{1}, \varphi_{2}, \ldots\right\}$ for $W_{0}$ as in Theorem 2.1. Since our tight framelet $\Psi$ consists of $N$ of functions, $\Phi$ has $\leqslant N$ of non-zero elements as well. Condition 4.1 assures that $\Phi$ is a semi-orthogonal wavelet. This allows us to calculate the spectral function of $V_{0}$ in terms of $\Phi$. Indeed, a formula from [15] gives us

$$
\sigma_{V_{0}}(\xi)=\sum_{\varphi \in \Phi} \sum_{j>0}\left|\hat{\varphi}\left(B^{j} \xi\right)\right|^{2}
$$

After integrating the above formula and using the fact that $\|\varphi\| \leqslant 1$ for all $\varphi \in \Phi$, we obtain that

$$
\int_{\mathbb{R}^{n}} \sigma_{V_{0}}=\sum_{\varphi \in \Phi}\|\varphi\|^{2} /(|\operatorname{det} A|-1) \leqslant N /(|\operatorname{det} A|-1) .
$$

Since $\int_{\mathbb{R}^{n}} \sigma_{V_{0}}=\int_{\mathbb{T}^{n}} m$, this shows that (D5) is satisfied.
In order to justify (D6) we use basic properties of the dimension function that are given in [15]. Since $V_{1}=V_{0} \oplus W_{0}$, we get that $\sum_{d \in \mathcal{D}} m\left(B^{-1} \xi+d\right)=m(\xi)+\operatorname{dim}_{W_{0}}(\xi)$. But $W_{0}$ has $N$ generators, therefore $\operatorname{dim}_{W_{0}} \leqslant N$ and (D6) follows.

Remark 4.1. In order to construct our GMRA only conditions (D1)-(D5) are going to be used. We shall also show that (D6) is necessary and sufficient to guarantee the existence of a "high pass filter" that will be used to define our framelet. We also want to point out, that (D5) follows from (D6), as was shown in Lemma 3.1.

In short, our construction is guided by the standard procedure. We are going to consider a matrix mask function $M$ that satisfies conditions (2.14)-(2.16). Then we will construct a corresponding refinable vector $\Phi$ and use Unitary Extension Principle to obtain an associated tight framelet $\Psi$.

We start equipped with a function $m$ that satisfies conditions (D1)-(D5). Then, we define the sets $S_{j}$, for $j \in \mathbb{N}$, by a formula analogous to (2.6), that is

$$
\begin{equation*}
S_{j}=\left\{\xi \in \mathbb{R}^{n}: m(\xi) \geqslant j\right\} \tag{4.2}
\end{equation*}
$$

Let $J=\left\{i \in \mathbb{N}:\left|S_{j}\right|>0\right\}$. Hence, $J=\{1,2, \ldots, L\}$ or $J=\mathbb{N}$. The sets $S_{j}, j \in J$, are used to define the diagonal matrix function $\Omega$ as in (2.9). This allows us to consider a matrix mask function $M$ with periodic entries $m_{i, j} \in L^{2}\left(S_{j}\right), i, j \in J$, that satisfies conditions (2.14)-(2.16). In order to find the corresponding refinable vector we shall use the ideas of [27]. First, we will modify $M$ to assure that the product of the dilates of $M$ is convergent. Then, we shall use multipliers to recover the solution to the original problem. To proceed in this direction we need the following basic lemma about multipliers.

Lemma 4.1. Let $\mu$ be a unimodular measurable function on $\mathbb{R}^{n}$ (that is, $\mu: \mathbb{R}^{n} \rightarrow S^{1}=$ $\{z \in \mathbb{C}:|z|=1\}$ ). If $B$ is an expansive matrix, then there exists a unimodular measurable function $v$ such that

$$
\begin{equation*}
\nu(B \xi) \overline{\nu(\xi)}=\mu(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{4.3}
\end{equation*}
$$

Proof. It is well known, that for any expansive matrix $B$ there is an ellipsoid $\mathcal{E}$ such that $\mathcal{E} \subset B(\mathcal{E})$, see e.g. [12, Lemma 2.2]. It follows, that for $W=B(\mathcal{E}) \backslash \mathcal{E}$ we have $\bigcup_{j \in \mathbb{Z}} B^{j}(W)=\mathbb{R}^{n}$. Therefore, it is enough to define a unimodular function $v$ on $W$ and then extend it to $\mathbb{R}^{n}$ using Eq. (4.3).

The mentioned modification of the matrix mask function $M$ is very simple. The most important entry of $M$ is $m_{1,1}$. Let $\mu$ be a phase of $m_{1,1}$. That is, $\mu$ is a unimodular measurable function such that

$$
\begin{equation*}
\mu(\xi)\left|m_{1,1}(\xi)\right|=m_{1,1}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{4.4}
\end{equation*}
$$

A multiplier associated to the mask $M$ is any unimodular measurable function $v$ satisfying (4.3) and (4.4). The modified mask is

$$
\begin{equation*}
M^{\prime}:=\bar{\mu} M \tag{4.5}
\end{equation*}
$$

and the corresponding refinable vector is given by

$$
\begin{equation*}
\hat{\Phi}^{\prime}(\xi):=\lim _{N \rightarrow \infty}\left[\prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)\right] e \tag{4.6}
\end{equation*}
$$

where $e$ is the vector $(1,0,0, \ldots)$. Finally, the refinable vector $\Phi$ corresponding to the mask $M$ will be given by $\hat{\Phi}:=\nu \hat{\Phi}^{\prime}$. In order to establish that $\Phi$ is refinable we use the following series of lemmas.

Lemma 4.2. The vector function $\hat{\Phi}^{\prime}$ in (4.6) is well defined.

Proof. We need to show, that the limit in (4.6) does exist for a.e. $\xi \in \mathbb{R}^{n}$. We want to point out, that although $\prod_{j=1}^{\infty} M^{\prime}\left(B^{-j} \xi\right)$ may not exist, we are only interested in the first column of this matrix. Later, we will prove that under some natural assumptions this product matrix exists and all of its columns but the first must be zero, see (5.1).

By (2.15), for a.e. $\xi \in \mathbb{R}^{n}$ we can find $N(\xi)$ such that the first column of $M^{\prime}\left(B^{-j} \xi\right)$ has only one non-zero entry (the first one) for all $j>N(\xi)$. Therefore, for $N>N(\xi)$ we have

$$
\left[\prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)\right] e=\left[\prod_{j=1}^{N(\xi)} M^{\prime}\left(B^{-j} \xi\right)\right]\left[\prod_{j=N(\xi)+1}^{N}\left|m_{1,1}\left(B^{-j} \xi\right)\right|\right] e .
$$

Thus,

$$
\hat{\Phi}^{\prime}(\xi):=\lim _{N \rightarrow \infty} p_{N}(\xi) v(\xi)
$$

where $v(\xi)=\left[\prod_{j=1}^{N(\xi)} M^{\prime}\left(B^{-j} \xi\right)\right] e$ and $p_{N}(\xi)=\prod_{j=N(\xi)+1}^{N}\left|m_{1,1}\left(B^{-j} \xi\right)\right|$. Since condition (2.14) guarantees that $\left|m_{1,1}\right| \leqslant 1$, we see that $\left\{p_{N}(\xi)\right\}$ is a bounded decreasing sequence and our claim follows.

Lemma 4.3. The vector function $\hat{\Phi}^{\prime}$ in (4.6) satisfies $\hat{\Phi}^{\prime}(B \xi)=M^{\prime}(\xi) \hat{\Phi}^{\prime}(\xi)$.
Proof. From (D5) it follows that our function $m$ is finite a.e. Therefore, condition (2.14) implies that for a.e. $\xi \in \mathbb{T}^{n}$ the matrix $M^{\prime}(\xi)$ has only finitely many non-zero terms. This allows us to see that

$$
\begin{aligned}
\hat{\Phi}^{\prime}(B \xi) & =\lim _{N \rightarrow \infty}\left(M^{\prime}(\xi)\left[\prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)\right] e\right)=M^{\prime}(\xi) \lim _{N \rightarrow \infty}\left(\left[\prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)\right] e\right) \\
& =M^{\prime}(\xi) \hat{\Phi}^{\prime}(\xi) .
\end{aligned}
$$

Lemma 4.4. The vector function $\hat{\Phi}^{\prime}$ in (4.6) satisfies $\lim _{N \rightarrow \infty}\left\|\hat{\Phi}^{\prime}\left(B^{-N} \xi\right)\right\|=1$, for a.e. $\xi \in \mathbb{R}^{n}$.
Proof. By (2.15), for a.e. $\xi \in \mathbb{R}^{n}$ we can find $N(\xi)$ such that for all $N>N(\xi)$ we have

$$
\hat{\Phi}^{\prime}\left(B^{-N} \xi\right)=\left(\prod_{j=N+1}^{\infty}\left|m_{1,1}\left(B^{-j} \xi\right)\right|\right) e
$$

Therefore, our claim follows from (2.16).
Lemma 4.5. The vector function $\hat{\Phi}^{\prime}$ in (4.6) satisfies $\int_{\mathbb{R}^{n}}\left\|\hat{\Phi}^{\prime}(\xi)\right\|^{2} d \xi<\infty$.
Proof. For $N \in \mathbb{N}$ and a.e. $\xi \in \mathbb{R}^{n}$ let us consider the following matrix

$$
\begin{equation*}
M_{N}(\xi)=\left[\prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)\right] \mathbf{1}_{B^{N}\left(\mathbb{T}^{n}\right)}(\xi) \tag{4.7}
\end{equation*}
$$

We claim that for all $N \in \mathbb{N}$ and a.e. $\xi \in \mathbb{R}^{n}$,

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}} M_{N}(\xi+k) M_{N}^{*}(\xi+k)=\Omega(\xi) \tag{4.8}
\end{equation*}
$$

Indeed, for $N=1$ we use (2.14) to obtain

$$
\begin{aligned}
\sum_{k \in \mathbb{Z}^{n}} M_{1}(\xi+k) M_{1}^{*}(\xi+k) & =\sum_{k \in \mathbb{Z}^{n}} M\left(B^{-1}(\xi+k)\right) M^{*}\left(B^{-1}(\xi+k)\right) \mathbf{1}_{B\left(\mathbb{T}^{n}\right)}(\xi+k) \\
& =\sum_{d \in \mathcal{D}} M\left(B^{-1} \xi+d\right) M^{*}\left(B^{-1} \xi+d\right)=\Omega(\xi)
\end{aligned}
$$

To proceed with the induction we observe that

$$
M_{N+1}(\xi)=M^{\prime}\left(B^{-1} \xi\right) M_{N}^{\prime}\left(B^{-1} \xi\right)
$$

for $N \in \mathbb{N}$ and a.e. $\xi \in \mathbb{R}^{n}$. Therefore,

$$
\begin{aligned}
& \sum_{k \in \mathbb{Z}^{n}} M_{N+1}(\xi+k) M_{N+1}^{*}(\xi+k) \\
& \quad=\sum_{k \in \mathbb{Z}^{n}} M\left(B^{-1}(\xi+k)\right) M_{N}\left(B^{-1}(\xi+k)\right) M_{N}^{*}\left(B^{-1}(\xi+k)\right) M^{*}\left(B^{-1}(\xi+k)\right) \\
& \quad=\sum_{d \in \mathcal{D}} \sum_{l \in \mathbb{Z}^{n}} M\left(B^{-1} \xi+d\right) M_{N}\left(B^{-1} \xi+d+l\right) M_{N}^{*}\left(B^{-1} \xi+d+l\right) M^{*}\left(B^{-1} \xi+d\right) \\
& \quad=\sum_{d \in \mathcal{D}} M\left(B^{-1} \xi+d\right) \Omega\left(B^{-1} \xi+d\right) M^{*}\left(B^{-1} \xi+d\right) \\
& \quad=\sum_{d \in \mathcal{D}} M\left(B^{-1} \xi+d\right) M^{*}\left(B^{-1} \xi+d\right)=\Omega(\xi)
\end{aligned}
$$

what proves our claim (4.8). In order to use it, we observe that for all $k \in \mathbb{Z}^{n}$ and a.e. $\xi \in \mathbb{R}^{n}$,

$$
\begin{equation*}
\left\|M_{N}(\xi+k) e\right\|^{2} \leqslant\left\|M_{N}(\xi+k)\right\|^{2} \leqslant\left\|M_{N}(\xi+k)\right\|_{H S}^{2}=\operatorname{tr}\left[M_{N}(\xi+k) M_{N}^{*}(\xi+k)\right] \tag{4.9}
\end{equation*}
$$

where $\|\cdot\|_{H S}$ denotes the Hilbert-Schmidt operator norm. The above estimate and (4.8) give us

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}}\left\|M_{N}(\xi+k) e\right\|^{2} \leqslant \operatorname{tr}[\Omega(\xi)]=m(\xi) \tag{4.10}
\end{equation*}
$$

Since $\lim _{N \rightarrow \infty}\left(M_{N}(\xi) e\right)=\hat{\Phi}^{\prime}(\xi)$, we can use Fatou's lemma to conclude that

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}}\left\|\hat{\Phi}^{\prime}(\xi+k)\right\|^{2} \leqslant m(\xi) \tag{4.11}
\end{equation*}
$$

By (D5) the function $m(\xi)$ is integrable over $\mathbb{T}^{n}$, thus

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left\|\hat{\Phi}^{\prime}(\xi)\right\|^{2} d \xi=\int_{\mathbb{T}^{n}} \sum_{k \in \mathbb{Z}^{n}}\left\|\hat{\Phi}^{\prime}(\xi+k)\right\|^{2} d \xi \leqslant \int_{\mathbb{T}^{n}} m(\xi) d \xi<\infty \tag{4.12}
\end{equation*}
$$

Remark 4.2. We have that $\hat{\Phi}^{\prime}=\left(\hat{\varphi}_{j}^{\prime}\right)_{j \in J}$. Since $\int_{\mathbb{R}^{n}}\left\|\hat{\Phi}^{\prime}(\xi)\right\|^{2} d \xi=\sum_{j \in J}\left\|\hat{\varphi}_{j}\right\|^{2}$, the above lemma shows that $\hat{\Phi}^{\prime} \subset L^{2}\left(\mathbb{R}^{n}\right)$.

To reverse the procedure given in (4.5) we use Lemma 4.1 to find a multiplier $v$ associated to $M$ and define our refinable vector $\Phi$ by setting

$$
\begin{equation*}
\hat{\Phi}:=v \hat{\Phi}^{\prime} \tag{4.13}
\end{equation*}
$$

The following result assures that such $\Phi$ has all of the properties that we need.
Theorem 4.2. The vector function $\hat{\Phi}$ given in (4.13) satisfies conditions (3.1)-(3.4).
Proof. By Lemma 4.3 and (4.3) together with (4.5) we get that

$$
\hat{\Phi}(B \xi)=v(B \xi) \hat{\Phi}^{\prime}(B \xi)=v(B \xi) M^{\prime}(\xi) \hat{\Phi}^{\prime}(\xi)=v(B \xi) \bar{\mu}(\xi) M(\xi) \bar{\nu}(\xi) \hat{\Phi}(\xi)=M(\xi) \hat{\Phi}(\xi)
$$

therefore, (3.1) holds. As we mentioned before, (2.14) implies that the mask matrix $M(\xi)$ has only finitely many non-zero terms. Thus, condition (3.2) is satisfied, by Remark 3.1. Since $\|\hat{\Phi}(\xi)\|=\left\|\hat{\Phi}^{\prime}(\xi)\right\|$ a.e., properties (3.3) and (3.4) follow immediately from Lemmas 4.4 and 4.5.

As an immediate consequence of Theorem 4.2 and the Unitary Extension Principle from the previous section, we obtain our framelet construction result. The precursor of Theorem 4.3 is a result of Baggett, Jorgensen, Merrill, and Packer [4, Theorem 3.4] where a low-pass matrix mask $M$ is assumed to be finite and Lipschitz continuous near 0 (instead of satisfying our assumptions (2.15) and (2.16)).

Theorem 4.3. Let $m$ be a function that satisfies (D1)-(D5) with the sets $S_{j}$ given by (4.2) and the corresponding matrix function $\Omega$ defined in (2.9). Let $M=\left(m_{i, j}\right)_{i, j \in J}$ be a matrix mask function with periodic entries $m_{i, j} \in L^{2}\left(S_{j}\right)$, that satisfies (2.14)-(2.16). Then there is a refinable vector $\Phi$ such that

$$
\begin{equation*}
\hat{\Phi}(B \xi)=M(\xi) \hat{\Phi}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{4.14}
\end{equation*}
$$

Moreover, if a matrix function $H=\left(h_{i, j}\right)_{i \in \tilde{J}, j \in J}$ with $\tilde{J}=\{1, \ldots, N\}$ finite and with periodic entries $h_{i, j} \in L^{2}\left(S_{j}\right)$ satisfies

$$
\begin{equation*}
M^{*}(\xi) M(\xi+d)+H^{*}(\xi) H(\xi+d)=\Omega(\xi) \delta_{0, d} \quad \text { for any } d \in \mathcal{D} \text { and a.e. } \xi \in \mathbb{R}^{n}, \tag{4.15}
\end{equation*}
$$

then $\Psi=\left(\psi^{j}\right)_{j \in \tilde{J}}$ given by

$$
\begin{equation*}
\hat{\Psi}(B \xi)=H(\xi) \hat{\Phi}(\xi) \tag{4.16}
\end{equation*}
$$

is a tight framelet for $L^{2}\left(\mathbb{R}^{n}\right)$.

Proof. The existence of a refinable vector $\Phi$ satisfying (4.14) is a consequence of Theorem 4.2. In addition, if a matrix function $H$ satisfies (4.15), then by Theorem 3.1, $\Psi$ given by (4.16) is a tight framelet for $L^{2}\left(\mathbb{R}^{n}\right)$.

The next theorem gives the necessary and sufficient conditions for the existence of the highpass matrix mask $H$ that satisfies (4.15). Some of the implications in Theorem 4.4 are already known. Indeed, (i) $\Rightarrow$ (ii) is due to Baggett, Courter, and Merrill [3, Theorem 2.5], whereas (ii) $\Rightarrow$ (iii) is due to Baggett, Jorgensen, Merrill, and Packer [4, Theorem 2.5] in the case of bounded $m$. We shall give the full proof of Theorem 4.4 for the sake of completeness.

Theorem 4.4. Let $m$ be any function satisfying (D1)-(D5). Let $\tilde{m}: \mathbb{R}^{n} \rightarrow \mathbb{N} \cup\{0\}$ be a measurable $\mathbb{Z}^{n}$-periodic function satisfying

$$
\begin{equation*}
\sum_{d \in \mathcal{D}} m(\xi+d)=m(B \xi)+\tilde{m}(B \xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{4.17}
\end{equation*}
$$

Define the sets

$$
S_{j}=\left\{\xi \in \mathbb{R}^{n}: m(\xi) \geqslant j\right\}, \quad \tilde{S}_{j}=\left\{\xi \in \mathbb{R}^{n}: \tilde{m}(\xi) \geqslant j\right\}
$$

and the corresponding matrix functions $\Omega$ and $\tilde{\Omega}$ by (2.9). Assume that $M=\left(m_{i, j}\right)_{i, j \in J}$ is a matrix mask function with periodic entries $m_{i, j} \in L^{2}\left(S_{j}\right)$ that satisfies (2.14). Then the following are equivalent:
(i) $m$ satisfies (D6).
(ii) There exists a matrix function $H=\left(h_{i, j}\right)_{i \in \tilde{J}, j \in J}$ with $\tilde{J}$ finite and with periodic entries $h_{i, j} \in L^{2}\left(S_{j}\right)$ satisfying

$$
\begin{gather*}
\sum_{d \in \mathcal{D}} H(\xi+d) H^{*}(\xi+d)=\tilde{\Omega}(B \xi) \quad \text { a.e. } \xi \in \mathbb{R}^{n}  \tag{4.18}\\
\sum_{d \in \mathcal{D}} H(\xi+d) M^{*}(\xi+d)=0 \quad \text { a.e. } \xi \in \mathbb{R}^{n} \tag{4.19}
\end{gather*}
$$

(iii) There exists a matrix function $H=\left(h_{i, j}\right)_{i \in \tilde{J}, j \in J}$ with $\tilde{J}$ finite and with periodic entries $h_{i, j} \in L^{2}\left(S_{j}\right)$ satisfying (4.15).

Moreover, if a matrix function $H$ satisfies (4.18) and (4.19) then it also satisfies (4.15). However, the converse is in general false.

Remark 4.3. Note that $\tilde{m}$ as in (4.17) always exists, since

$$
\tilde{m}(\xi)=\sum_{d \in \mathcal{D}} m\left(B^{-1} \xi+d\right)-m(\xi)
$$

is clearly $\mathbb{Z}^{n}$-periodic and non-negative by (D2). Moreover, let $E$ be any measurable subset of $\mathbb{T}^{n}$ such that $\{E+d: d \in \mathcal{D}\}$ is a partition of $\mathbb{T}^{n}$ (modulo null sets). Then, it is easy to see using the
periodicity of $M$ and $H$ that if (4.18) and (4.19) hold for a.e. $\xi \in E$, then they must hold for a.e. $\xi \in \mathbb{R}^{n}$.

Proof. First, suppose that a matrix function $H=\left(h_{i, j}\right)_{i \in \tilde{J}, j \in J}$ has periodic entries $h_{i, j} \in$ $L^{2}\left(S_{j}\right)$, satisfies (4.15) and the index set $\tilde{J}$ has $N$ elements. Consider the combined matrix function

$$
T(\xi)=\left[\begin{array}{lll}
M\left(\xi+d_{1}\right) & \ldots & M\left(\xi+d_{q}\right)  \tag{4.20}\\
H\left(\xi+d_{1}\right) & \ldots & H\left(\xi+d_{q}\right)
\end{array}\right]
$$

By the support conditions and (4.15), the matrix $T(\xi)$ has precisely $\sum_{d \in \mathcal{D}} m(\xi+d)$ non-zero columns. Condition (4.15) says that these non-zero columns form an orthonormal system. On the other hand, (2.14) and the fact that $\tilde{J}$ has $N$ elements imply that the matrix $T(\xi)$ has at most $m(B \xi)+N$ non-zero rows. Clearly, any collection of orthonormal vectors must be smaller than the dimension of the space where they live in. Consequently, (D6) must hold. This shows (iii) $\Rightarrow$ (i).

Conversely, suppose that (D6) holds and consider a matrix function

$$
T^{\prime}(\xi)=\left[\begin{array}{lll}
M\left(\xi+d_{1}\right) & \ldots & M\left(\xi+d_{q}\right) \tag{4.21}
\end{array}\right]
$$

It is convenient to fix $\xi \in E$, where $E$ is the same as in Remark 4.3. As before, by the support conditions, the matrix $T^{\prime}(\xi)$ has at most $\sum_{d \in \mathcal{D}} m(\xi+d)$ non-zero columns. On the other hand, the matrix $T^{\prime}(\xi)$ has $m(B \xi)$ non-zero rows forming an orthonormal system by (2.14). Therefore, for a fixed $\xi$, we have a finite submatrix $T^{\prime \prime}(\xi)$ with $c=\sum_{d \in \mathcal{D}} m(\xi+d)$ columns and $r=$ $m(B \xi)$ orthonormal rows. Now, it suffices to find an extension of this submatrix to a unitary $c \times c$ matrix. Since $\tilde{m}(B \xi)=c-r \leqslant N$ by (D6), at most $N$ extra rows must be added. Define $\left[H\left(\xi+d_{1}\right) \ldots H\left(\xi+d_{q}\right)\right]$ to be a matrix with rows indexed by $\tilde{J}=\{1, \ldots, N\}$ such that the first $\tilde{m}(B \xi)=c-r$ rows of $\left[H\left(\xi+d_{1}\right) \ldots H\left(\xi+d_{q}\right)\right]$ are formed by inserting the extra rows from a finite submatrix $T^{\prime \prime}(\xi)$ interspersed by zero columns, which were previously removed from the matrix $T^{\prime}(\xi)$. The remaining rows (if any) of $\left[H\left(\xi+d_{1}\right) \ldots H\left(\xi+d_{q}\right)\right]$ are defined to be zero. It is not hard to see that the these extra rows can be chosen in such a way that the resulting matrix function $\left[H\left(\xi+d_{1}\right) \ldots H\left(\xi+d_{q}\right)\right]$ has measurable entries as a function of $\xi \in E$. As a result, the combined matrix function $T(\xi)$ has the same number of non-zero columns equal to $\sum_{d \in \mathcal{D}} m(\xi+d)$ as the number of non-zero rows equal to $m(B \xi)+\tilde{m}(B \xi)$ by (4.17). Furthermore, since the non-zero rows of $T(\xi)$ form an orthonormal sequence, the finite submatrix consisting of non-zero columns and rows must be unitary. Consequently, the constructed matrix $H$ satisfies (4.18) and (4.19) for a.e. $\xi \in E$. By Remark 4.3 this shows that (i) $\Rightarrow$ (ii).

Next, if $H$ is any matrix function as in (ii), then (2.14), (4.18), and (4.19) imply that the non-zero rows of the combined matrix function $T(\xi)$ form an orthonormal sequence. By (4.17) a finite submatrix consisting of non-zero columns and rows of $T(\xi)$ has the same number of non-zero columns as the number of non-zero rows and hence must be unitary. Since the rows of this finite submatrix are orthonormal, so are the columns, which implies that (4.15) holds. Therefore, (4.18) and (4.19) always imply (4.15). The converse implication is obviously false in general, since the combined matrix $T(\xi)$ may have a larger number of non-zero rows than non-zero columns and as a consequence the orthonormality of columns does not translate into orthonormality of rows. This proves (ii) $\Rightarrow$ (iii) and completes the proof of Theorem 4.4.

Remark 4.4. The origin of Eq. (4.17) is hidden in the consistency equation (2.22). Once we take $m=\operatorname{dim}_{V_{0}}$ and $\tilde{m}=\operatorname{dim}_{W_{0}}$, the connection becomes clear.

As a consequence of Theorems 4.3 and 4.4 we can deduce that any low-pass matrix mask function $M$ satisfying (2.14)-(2.16) associated with the dimension function $m$ satisfying (D1)(D6) corresponds to some tight framelet $\Psi$ via (4.14) and (4.16). To achieve this we must choose a high-pass mask matrix $H$ such that the corresponding combined matrix function (4.20) has orthogonal columns, that is, (4.15) holds. Naturally, if we count on obtaining a semi-orthogonal wavelet $\Psi$, then the high-pass matrix mask $H$ must satisfy more restrictive conditions (4.18)(4.19) resulting in row orthogonality of the combined matrix (4.20), see Theorem 5.2.

We would like to point out, that a refinable vector $\Phi$ in Theorem 4.3 is not unique. Indeed, the explicit formula for our choice of $\Phi$ is

$$
\begin{equation*}
\hat{\Phi}(\xi)=v(\xi) \lim _{N \rightarrow \infty}\left[\prod_{j=1}^{N} \bar{\mu}\left(B^{-j} \xi\right) M\left(B^{-j} \xi\right)\right] e=\lim _{N \rightarrow \infty} v\left(B^{-N} \xi\right)\left[\prod_{j=1}^{N} M\left(B^{-j} \xi\right)\right] e \tag{4.22}
\end{equation*}
$$

where $e=(1,0,0, \ldots), \mu$ is the phase of $m_{1,1}$ and $\nu$ is an arbitrary multiplier, i.e., a measurable unimodular function such that $\nu(B \xi) \overline{\nu(\xi)}=\mu(\xi)$. Recall that Lemma 4.1 guarantees the existence of such multipliers. Equivalently, we can define a multiplier associated to the mask $M=\left(m_{i, j}\right)$ as any function $v$ satisfying

$$
\begin{equation*}
v(B \xi) \overline{v(\xi)}\left|m_{1,1,}(\xi)\right|=m_{1,1}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} . \tag{4.23}
\end{equation*}
$$

Indeed, if $v$ satisfies (4.23), then $\mu(\xi)=v(B \xi) \overline{v(\xi)}$ is a phase of $m_{1,1}$ and $v$ is its corresponding multiplier. Note that a phase $\mu$ satisfying (4.4) might not be unique if $m_{1,1}$ does not have a full support.

Since there are many possibilities for multipliers $v$ satisfying (4.23) we obtain a lot of choices for $\Phi$. Moreover, these different choices generate distinct GMRA's. In general, if $P$ is a matrix function such that $P(B \xi)^{-1} M(\xi) P(\xi)=M(\xi)$ a.e., then our $\hat{\Phi}$ can be replaced by $P \hat{\Phi}$. Nevertheless, we can loosely think that $\Phi$ is given by the standard product $\prod_{j=1}^{\infty} M\left(B^{-j} \xi\right)$ that is applied to the vector $e$. Even better, it turns out that if the product is convergent, then this standard choice of $\Phi$ is valid.

Proposition 4.2. If $M$ is as in Theorem 4.3 and the product $\prod_{j=1}^{\infty} M\left(B^{-j} \xi\right)$ is convergent for a.e. $\xi \in \mathbb{R}^{n}$, then $\Phi$ from Theorem 4.3 can be taken as

$$
\begin{equation*}
\hat{\Phi}(\xi):=\left[\prod_{j=1}^{\infty} M\left(B^{-j} \xi\right)\right] e . \tag{4.24}
\end{equation*}
$$

Proof. Since $\prod_{j=1}^{\infty} M\left(B^{-j} \xi\right)$ is convergent a.e., we have that $\prod_{j=1}^{\infty} m_{1,1}\left(B^{-j} \xi\right)$ is convergent a.e. as well. In particular,

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \prod_{j=N}^{\infty} m_{1,1}\left(B^{-j} \xi\right)=1 \quad \text { a.e. } \tag{4.25}
\end{equation*}
$$

By (4.3) this is equivalent to

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \nu\left(B^{-N} \xi\right) \quad \text { exists for a.e. } \xi \in \mathbb{R}^{n} \tag{4.26}
\end{equation*}
$$

Let $\Phi$ be a refinable vector as in Theorem 4.3 and $v$ be any function as in Lemma 4.1. Define another function $\nu^{\prime}$ satisfying the conclusions of Lemma 4.1 by

$$
\nu^{\prime}(\xi)=\overline{\alpha(\xi)} \nu(\xi), \quad \text { where } \alpha(\xi)=\lim _{N \rightarrow \infty} v\left(B^{-N} \xi\right)
$$

Indeed,

$$
v^{\prime}(B \xi) \overline{v^{\prime}(\xi)}=\mu(\xi), \quad \lim _{N \rightarrow \infty} v^{\prime}\left(B^{-N} \xi\right)=1 \quad \text { for a.e. } \xi \in \mathbb{R}^{n}
$$

Therefore, by (4.22),

$$
\hat{\Phi}^{\prime}(\xi)=\left[\prod_{j=1}^{\infty} M\left(B^{-j} \xi\right)\right] e
$$

is a refinable vector function obtained by the procedure of Theorem 4.3 with the multiplier $v^{\prime}$. This proves (4.24).

## 5. Reconstruction of wavelets

The main goal of this section is to prove that every orthogonal wavelet can be reconstructed from its carefully chosen low-pass and high-pass matrix masks by the procedure described in the previous section. In order to achieve this, we will explore in more depth some subtle properties of the refinable vector $\Phi$ from Theorem 4.3.

Recall that by starting from a dimension function $m$ and an appropriate matrix mask $M$, we obtained a refinable vector $\Phi$ in Section 4 and, therefore, also the associated GMRA. The standard issue in this type of constructions is the problem of "vanishing mass." In short, it may happen that $\int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2} d \xi<\int_{\mathbb{R}^{n}} m(\xi) d \xi$. In particular, $\Phi$ need not to be a scaling vector since quasi-orthogonality may fail. Also, the GMRA that results from such procedure can have a strictly smaller dimension function (of its core space) than the original one that was used to start the construction. This feature was already observed in the classical MRA case on $\mathbb{R}$ with dilation by 2. The familiar Cohen's condition is one of the ways to assure that "no mass gets lost." It is crucial if one hopes to obtain a wavelet. However, as pointed in [27] in the dyadic scalar case, even if "some of the mass does vanish" one can still construct corresponding tight framelet. In the general case, the problem gains on complexity. Below, we give a simple necessary condition that is needed for preserving the "mass."

In order to achieve this preservation, one has to impose that the matrix $M^{\prime}$ given in (4.5) satisfies

$$
\lim _{N \rightarrow \infty} \prod_{j=1}^{N} M^{\prime}\left(B^{-j} \xi\right)=\left[\begin{array}{cccc}
* & 0 & 0 & \ldots  \tag{5.1}\\
* & 0 & 0 & \ldots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

More precisely, for a.e. $\xi \in \mathbb{R}^{n}$ the limit does exist and is equal to a matrix whose only non-zero entries are in the first column. This also sheds a new light on (4.6), where we defined $\hat{\Phi}^{\prime}$ as the first column of such a product.

We will show the necessity of the above condition in the following
Proposition 5.1. Let $m$ and $\Phi$ be as in Theorem 4.3. If

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2} d \xi=\int_{\mathbb{T}^{n}} m(\xi) d \xi \tag{5.2}
\end{equation*}
$$

then (5.1) holds.
Proof. Let us show (5.1) by using calculations given in the proof of Lemma 4.5. First, we observe that (5.1) is equivalent to saying that $\lim _{N \rightarrow \infty}\left\|M_{N}(\xi) e_{i}\right\| \rightarrow 0$ for a.e. $\xi \in \mathbb{R}^{n}$, and every vector $e_{i}, i \geqslant 2$, of the standard basis. Also, we can use $\Phi^{\prime}$ instead of $\Phi$ in our considerations.

By (4.12), the assumption (5.2) forces the inequality (4.11) to become an equality. Tracing back this fact through (4.10) and (4.9) we see that, eventually, one must have

$$
\lim _{N \rightarrow \infty}\left(\operatorname{tr}\left[M_{N}(\xi) M_{N}^{*}(\xi)\right]-\left\|M_{N}(\xi) e\right\|^{2}\right)=0
$$

for a.e. $\xi \in \mathbb{R}^{n}$. However, since $\operatorname{tr}[C D]=\operatorname{tr}[D C]$, the above becomes

$$
\lim _{N \rightarrow \infty} \sum_{i \geqslant 2}\left\|M_{N}(\xi) e_{i}\right\|^{2}=0
$$

This shows the necessity of (5.1) and concludes the proof.
In the next result we present the full connection between the "mass preservation" and the properties of our refinable vector $\Phi$.

Theorem 5.1. Let $m$ and $\Phi$ be as in Theorem 4.3. Then, $\Phi$ is a scaling vector that generates $a$ GMRA with the same dimension function as $m$ if and only if (5.2) holds.

Proof. Again, we can consider $\Phi^{\prime}$ instead of $\Phi$. If $\Phi^{\prime}$ is a scaling vector then the dimension function of the corresponding GMRA is equal to $\sum_{k \in \mathbb{Z}^{n}}\left\|\hat{\Phi}^{\prime}(\xi+k)\right\|^{2}$. Clearly, the assumption that this dimension function is the same as $m$ implies that (5.2) holds.

On the other hand, assume that (5.2) is satisfied. From (4.12) and (4.11) it follows that

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}}\left\|\hat{\Phi}^{\prime}(\xi+k)\right\|^{2}=m(\xi) \tag{5.3}
\end{equation*}
$$

for a.e. $\xi \in \mathbb{T}^{n}$. By Proposition 5.1, applying Fatou's lemma to (4.8) yields

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}^{n}} \hat{\Phi}^{\prime}(\xi+k) \hat{\Phi}^{\prime *}(\xi+k) \leqslant \Omega(\xi) \tag{5.4}
\end{equation*}
$$

in the operator sense, for a.e. $\xi \in \mathbb{T}^{n}$. However, condition (5.3) simply says that

$$
\operatorname{tr}\left[\sum_{k \in \mathbb{Z}^{n}} \hat{\Phi}^{\prime}(\xi+k) \hat{\Phi}^{\prime *}(\xi+k)\right]=\operatorname{tr}[\Omega(\xi)] \quad \text { a.e. }
$$

Therefore, we must have an equality in (5.4). This shows that $\Phi^{\prime}$ is a scaling vector. Moreover, (5.3) assures that the GMRA generated by $\Phi^{\prime}$ has the dimension function equal to $m$.

As a consequence of Theorem 5.1 we show that the procedure of Theorem 4.3 can result in a semi-orthogonal wavelet (with the expected size of generators) only if the combined matrix (4.20) has orthogonal rows. As a corollary, we conclude that the necessary condition for constructing orthogonal wavelets is that the high-pass filter $H$ satisfies (4.18) and (4.19) with the diagonal matrix function $\tilde{\Omega}$ constantly equal to the identity matrix.

Theorem 5.2. Suppose that $m$ is a function that satisfies (D1)-(D6) and $\tilde{m}$ is given by (4.17). Suppose that $M$ and $H$ are low-pass and high-pass matrix masks as in Theorem 4.3. Let $\Psi$ be the corresponding tight framelet. Then

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\|\hat{\Psi}(\xi)\|^{2} d \xi \leqslant \int_{\mathbb{T}^{n}} \tilde{m}(\xi) d \xi \tag{5.5}
\end{equation*}
$$

Moreover, if $\Psi$ is a semi-orthogonal wavelet such that the equality holds in (5.5), then the highpass filter $H$ necessarily satisfies (4.18) and (4.19) with the diagonal matrix function $\tilde{\Omega}$ given by (2.19).

Proof. Let $\Phi$ be the refinable vector constructed in Theorem 4.3. Recall that the proof of Theorem 3.1 yields

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\|\hat{\Psi}(\xi)\|^{2} d \xi=(|\operatorname{det} A|-1) \int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2} d \xi<\infty \tag{5.6}
\end{equation*}
$$

On the other hand, since (D6) holds, condition (4.17) implies that $\tilde{m}$ is bounded. Therefore, we can integrate (4.17) over $\mathbb{T}^{n}$ to obtain

$$
\begin{equation*}
\int_{\mathbb{T}^{n}} \tilde{m}(\xi) d \xi=(|\operatorname{det} A|-1) \int_{\mathbb{T}^{n}} m(\xi) d \xi<\infty \tag{5.7}
\end{equation*}
$$

Combining (4.12), (5.6), and (5.7) yields (5.5).
In addition, suppose that $\Psi$ is a semi-orthogonal wavelet such that the equality holds in (5.5). By Theorem 5.1, $\Phi$ is a scaling vector generating a GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ with the dimension function $\operatorname{dim}_{V_{0}}=m$. On the other hand, $\Psi$ also generates a GMRA $\left\{V_{j}^{\prime}\right\}_{j \in \mathbb{Z}}$ given by

$$
V_{j}^{\prime}=\operatorname{span}\left\{D^{i} T_{k} \psi^{\tilde{j}}: i<j, k \in \mathbb{Z}^{n}, \tilde{j} \in \tilde{J}\right\} \quad \text { for } j \in \mathbb{Z}
$$

By [15, Corollary 4.3] the dimension function $\operatorname{dim}_{V_{0}^{\prime}}$ of the core space $V_{0}^{\prime}$ can be computed explicitly and equals the wavelet dimension function $D_{\Psi}$. Consequently,

$$
\begin{equation*}
\int_{\mathbb{T}^{n}} \operatorname{dim}_{V_{0}^{\prime}}(\xi) d \xi=\frac{1}{|\operatorname{det} A|-1} \int_{\mathbb{R}^{n}}\|\hat{\Psi}(\xi)\|^{2}=\int_{\mathbb{R}^{n}}\|\hat{\Phi}(\xi)\|^{2}=\int_{\mathbb{T}^{n}} \operatorname{dim}_{V_{0}}(\xi) d \xi<\infty \tag{5.8}
\end{equation*}
$$

On the other hand, by (4.16) $\Psi \subset V_{1}$ and hence $V_{0}^{\prime} \subset V_{0}$. Thus, $\operatorname{dim}_{V_{0}^{\prime}}(\xi) \leqslant \operatorname{dim}_{V_{0}}(\xi)$ for a.e. $\xi$ and (5.8) implies that

$$
\operatorname{dim}_{V_{0}^{\prime}}(\xi)=\operatorname{dim}_{V_{0}}(\xi)<\infty \quad \text { for a.e. } \xi
$$

and hence $V_{0}^{\prime}=V_{0}$. Therefore, the semi-orthogonal wavelet $\psi$ is associated with the GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$. By Proposition 2.2 the high-pass matrix mask $H$ satisfies claimed properties.

As an immediate consequence of Theorem 5.2 we have
Corollary 5.1. In addition to the assumptions of Theorem 5.2, assume that the equality holds in (D6). Hence, there is $N \in \mathbb{N}$ such that

$$
\sum_{d \in \mathcal{D}} m(\xi+d)=m(B \xi)+N \quad \text { for a.e. } \xi \in \mathbb{R}^{n}
$$

Let $\Psi$ be the tight framelet as in Theorem 4.3. If $\Psi=\left\{\psi^{1}, \ldots, \psi^{N}\right\}$ is an orthonormal wavelet, then the high-pass filter $H$ necessarily satisfies (4.18) and (4.19) with the diagonal matrix function $\tilde{\Omega}(\xi) \equiv I d_{N \times N}$.

Proof. Our hypotheses imply that $\tilde{m}(\xi) \equiv N$ a.e. $\xi \in \mathbb{T}^{n}$. If $\Psi=\left(\psi^{j}\right)_{j=1}^{N}$ is an orthonormal wavelet, then

$$
\int_{\mathbb{R}^{n}}\|\hat{\Psi}(\xi)\|^{2} d \xi=N=\int_{\mathbb{T}^{n}} \tilde{m}(\xi) d \xi
$$

By Theorem 5.2 and (2.19) $H$ satisfies claimed properties since $\tilde{S}_{j}=\operatorname{supp} \operatorname{dim}_{\mathcal{S}\left(\psi^{j}\right)}=\mathbb{R}^{n}$ for $j=1, \ldots, N$.

Finally, we will show that every exhausting scaling vector of a GMRA (that has an integrable dimension function of the core space) can be obtained by the procedure of Theorem 4.3 with an appropriate choice of a multiplier $v$.

Theorem 5.3. Suppose $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is a GMRA such that $\operatorname{dim}_{V_{0}} \in L^{1}\left(\mathbb{T}^{n}\right)$. Let $M$ be the low-pass matrix mask function of an exhausting scaling vector $\Phi$ for the core space $V_{0}$.
(i) Any multiplier $v^{\prime}$ associated to $M$ corresponds by Theorem 4.3 to some scaling vector $\Phi^{\prime}$ with the same mask $M$ (but generating not necessarily the same space $V_{0}$ ).
(ii) There exists a multiplier $v$ associated to $M$ such that the scaling vector $\Phi$ is recovered via Theorem 4.3 by the product formula (4.22), that is

$$
\begin{equation*}
\hat{\Phi}(\xi)=\lim _{N \rightarrow \infty} v\left(B^{-N} \xi\right)\left[\prod_{j=1}^{N} M\left(B^{-j} \xi\right)\right] e \quad \text { for a.e. } \xi . \tag{5.9}
\end{equation*}
$$

Proof. Note that the dimension function $m=\operatorname{dim}_{V_{0}}$ of the core space $V_{0}$ satisfies the assumptions (D1)-(D5). Moreover, by Theorem 2.2, the matrix mask function $M$ of $\Phi$ satisfies conditions (2.14)-(2.16). Therefore, a fixed multiplier $v^{\prime}$ produces a refinable vector $\Phi^{\prime}$ by the procedure of Theorem 4.3.

We are going to prove (i) and (ii) simultaneously. Observe that both $\Phi$ and $\Phi^{\prime}$ satisfy the same refinable equation, which takes the form

$$
\begin{align*}
\hat{\Phi}\left(B^{-N} \xi\right) & =m_{1,1}\left(B^{-N-1} \xi\right) \hat{\Phi}\left(B^{-N-1} \xi\right) \\
\hat{\Phi}^{\prime}\left(B^{-N} \xi\right) & =m_{1,1}\left(B^{-N-1} \xi\right) \hat{\Phi}^{\prime}\left(B^{-N-1} \xi\right) \tag{5.10}
\end{align*}
$$

for sufficiently large $N>N(\xi)$ dependent on the choice of $\xi \in \mathbb{R}^{n}$. This is a simple consequence of the special form of the matrix mask $M$ near the origin. Let $\hat{\varphi}_{1}$ and $\hat{\varphi}_{1}^{\prime}$ be the first entries of $\hat{\Phi}$ and $\hat{\Phi}^{\prime}$, respectively. By (5.10), the sequence $\left\{\hat{\varphi}_{1}^{\prime}\left(B^{-N} \xi\right) / \hat{\varphi}_{1}\left(B^{-N} \xi\right)\right\}_{N>N(\xi)}$ must be constant whenever it is well defined. Let $\alpha(\xi)$ be the constant value of this sequence. It is clear that $\alpha(B \xi)=\alpha(\xi)$. Moreover, the fact that $\hat{\Phi}$ and $\hat{\Phi}^{\prime}$ have zeros in all but the first entry near the origin, (2.17), and Lemma 4.4, imply that $|\alpha(\xi)|=1$.

Define another multiplier $v$ corresponding to the same matrix mask function $M$ by

$$
v(\xi)=\overline{\alpha(\xi)} v^{\prime}(\xi)
$$

Finally, let $\Phi^{\prime \prime}$ be the refinable vector obtained by the procedure of Theorem 4.3 with the multiplier $\nu$. By (5.10) and the previously mentioned special form of $\Phi$ and $\Phi^{\prime}$ near the origin we have $\varepsilon>0$ such that

$$
\hat{\Phi}^{\prime}(\xi)=\alpha(\xi) \hat{\Phi}(\xi) \quad \text { for a.e. }|\xi|<\varepsilon
$$

On the other hand, by the product formula (4.22) we have that

$$
\begin{equation*}
\hat{\Phi}^{\prime}(\xi)=\alpha(\xi) \hat{\Phi}^{\prime \prime}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n} \tag{5.11}
\end{equation*}
$$

Since all functions $\Phi, \Phi^{\prime}$, and $\Phi^{\prime \prime}$ satisfy the refinable equation with respect to the same matrix mask function $M$, we must necessarily have that $\Phi=\Phi^{\prime \prime}$. This completes the proof of part (ii).

To deduce part (i) observe that $\Phi=\Phi^{\prime \prime}$ together with (5.11) yields

$$
\hat{\Phi}^{\prime}(\xi)=\alpha(\xi) \hat{\Phi}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R}^{n},
$$

for some $\alpha$ such that $|\alpha(\xi)|=1$ and $\alpha(B \xi)=\alpha(\xi)$ a.e. Therefore, we conclude that the refinable vector $\Phi^{\prime}$ must be necessarily a scaling vector. Obviously, there is no guarantee that the SI
space $V_{0}^{\prime}$ generated by $\Phi^{\prime}$ coincides with $V_{0}$ unless function $\alpha$ is $\mathbb{Z}^{n}$-periodic. This completes the proof of part (i) of Theorem 5.3.

As an immediate consequence of Theorem 5.3 we have that every semi-orthogonal wavelet $\Psi$ can be recovered by the procedure of Theorem 4.3. This is our main wavelet reconstruction result.

Theorem 5.4. Suppose $\Psi$ is a semi-orthogonal wavelet. Let $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ be the GMRA generated by $\Psi$, and let $\Phi$ be an exhausting scaling vector for $V_{0}$. Let $M$ and $H$ be the matrix mask functions of $\Phi$ and $\Psi$, respectively. Then there exists a multiplier $v$ associated to $M$ such that the scaling vector $\Phi$ is recovered by the product formula (5.9) and $\Psi$ is recovered by (4.16).

Proof. Theorem 5.3 guarantees that we can recover $\Phi$. To get back $\Psi$ we use Proposition 2.2. It implies that the matrix mask functions $M$ and $H$ satisfy (2.20) and (2.21). As we pointed out in Theorem 4.4, these two conditions force $H$ to satisfy (4.15). Thus, we can use such $H$ to obtain $\Psi$ via (4.16).

## 6. Examples and comments

We shall construct examples of tight framelets and wavelets using the procedure that was described in Section 4. We remind the reader that the whole process starts from choosing a dimension function. In order to find a specific dimension function one can construct a wavelet set and calculate the associated dimension function. In this way all dimension functions can be obtained by the result of Speegle and the authors [17].

It is customary to test GMRA constructions on the original "non-MRA object," that is, the Journé wavelet $\psi$ given by $\hat{\psi}=\mathbf{1}_{W}$, where $W=\left[-\frac{16}{7},-2\right] \cup\left[-\frac{1}{2},-\frac{2}{7}\right] \cup\left[\frac{2}{7}, \frac{1}{2}\right] \cup\left[2, \frac{16}{7}\right]$. The associated dimension function, called here the Journé dimension function, is

$$
m(\xi)= \begin{cases}2 & \text { for } \xi \in\left[-\frac{1}{7}, \frac{1}{7}\right]  \tag{6.1}\\ 1 & \text { for } \xi \in\left[-\frac{1}{2},-\frac{3}{7}\right] \cup\left[-\frac{2}{7},-\frac{1}{7}\right] \cup\left[\frac{1}{7}, \frac{2}{7}\right] \cup\left[\frac{3}{7}, \frac{1}{2}\right], \\ 0 & \text { for } \xi \in\left[-\frac{3}{7},-\frac{2}{7}\right] \cup\left[\frac{2}{7}, \frac{3}{7}\right]\end{cases}
$$

Since $m$ is $\mathbb{Z}$-periodic we list only its values on the torus $\mathbb{T}$ identified with $[-1 / 2,1 / 2)$. The same convention shall be used for other periodic objects that appear in this section albeit with the identification $\mathbb{T}=[-3 / 7,4 / 7)$. Clearly, $m$ satisfies conditions (D1)-(D6) that are stated in Theorem 4.1 and thereafter. The corresponding sets $S_{j}$ of (4.2) are $S_{1}$ and $S_{2}$, where $S_{1}$ is the periodization of $\left[-\frac{1}{2},-\frac{3}{7}\right] \cup\left[-\frac{2}{7}, \frac{2}{7}\right] \cup\left[\frac{3}{7}, \frac{1}{2}\right]$ and $S_{2}$ is the periodization of $\left[-\frac{1}{7}, \frac{1}{7}\right]$. The diagonal matrix function $\Omega$ of (2.9) is, therefore,

$$
\Omega(\xi)=\left[\begin{array}{cc}
\left.\mathbf{1}_{\left[-\frac{1}{2},-\frac{3}{7}\right]}\right]\left[-\frac{2}{7}, \frac{2}{7}\right] \cup\left[\frac{3}{7}, \frac{1}{2}\right] \\
0 & 0 \\
& \mathbf{1}_{\left[-\frac{1}{7}, \frac{1}{7}\right]}(\xi)
\end{array}\right]
$$

for $\xi \in \mathbb{T}$. All of this provides the ground for constructing tight framelets and wavelets by choosing appropriate low-pass and high-pass matrix mask functions.

We should mention that there are already several interesting constructions of wavelets and tight framelets with the Journé dimension function in the literature. For example, Baggett, Courter, and Merrill [3] constructed an orthonormal wavelet $\psi$ with the dimension function (6.1)
such that $\hat{\psi}$ is $C^{\infty}$ on an arbitrarily large interval. Later, Baggett, Jorgensen, Merrill, and Packer [4,5] gave an impressive construction of a tight framelet $\psi$ with a prescribed smoothness $\psi \in C^{r}$ for $r \geqslant 0$, and a global smoothness in the frequency $\hat{\psi} \in C^{\infty}$.

Our goal is to present another construction resulting in a large class of non-MSF wavelets sharing the same Journé dimension function. Unlike the previous constructions in [3-5], we shall insist that the first row of a low-pass matrix mask $M$ remains unchanged compared with that of the Journé wavelet.

Example 6.1. Consider the low-pass matrix mask function $M$ given by

$$
M(\xi)=\left[\begin{array}{cc}
\mathbf{1}_{F_{1}}(\xi) & 0  \tag{6.2}\\
m_{1}(\xi) & m_{2}(\xi)
\end{array}\right],
$$

where $F_{1}$ is the periodization of $\left[-\frac{2}{7},-\frac{1}{4}\right] \cup\left[-\frac{1}{7}, \frac{1}{7}\right] \cup\left[\frac{1}{4}, \frac{2}{7}\right]$ and $m_{1}, m_{2}$ are $\mathbb{Z}$-periodic measurable functions. Condition (2.14) imposes certain restrictions on possible functions $m_{1}$ and $m_{2}$. That is, we must stipulate that for $\xi \in \mathbb{T}$,

$$
\begin{align*}
\mathbf{1}_{F_{1}}(\xi) m_{1}(\xi)+\mathbf{1}_{F_{1}}(\xi+1 / 2) m_{1}(\xi+1 / 2) & =0  \tag{6.3}\\
\left|m_{1}(\xi)\right|^{2}+\left|m_{2}(\xi)\right|^{2}+\left|m_{1}(\xi+1 / 2)\right|^{2}+\left|m_{2}(\xi+1 / 2)\right|^{2} & =\mathbf{1}_{\left[-\frac{1}{14}, \frac{1}{14}\right] \cup\left[\frac{3}{7}, \frac{4}{7}\right]}(\xi) . \tag{6.4}
\end{align*}
$$

Here and subsequently, we are using the identification $\mathbb{T}=[-3 / 7,4 / 7)$. Since $F_{1}$ and $F_{1}+1 / 2$ are disjoint, $m_{1}$ must vanish on $F_{1}$ by (6.3) and consequently $m_{1}$ must be supported on the periodization of the interval $\left[-\frac{1}{14}, \frac{1}{14}\right]+1 / 2=\left[\frac{3}{7}, \frac{4}{7}\right]$ by (6.4). Consequently, we must have for $\xi \in \mathbb{T}$,

$$
\begin{equation*}
m_{1}(\xi)=v(\xi) \mathbf{1}_{\left[\frac{3}{7}, \frac{4}{7}\right]}(\xi), \quad m_{2}(\xi)=v(\xi) \mathbf{1}_{\left[-\frac{1}{14}, \frac{1}{14}\right]}(\xi) \tag{6.5}
\end{equation*}
$$

where $v$ is an arbitrary $\mathbb{Z}$-periodic measurable function satisfying

$$
\begin{equation*}
|v(\xi)|^{2}+|v(\xi+1 / 2)|^{2}=\mathbf{1}_{\left[-\frac{1}{14}, \frac{1}{14}\right]}(\xi), \quad \text { for a.e. } \xi \in(-1 / 4,1 / 4) \tag{6.6}
\end{equation*}
$$

It is easy to verify that as long as conditions (6.5) and (6.6) hold, the matrix mask function $M$ satisfies (2.14)-(2.16). Thus, we can apply Theorem 4.3. The corresponding refinable vector $\hat{\Phi}$ is the first column of the infinite product

$$
\prod_{j=1}^{\infty} M\left(2^{-j} \xi\right)=\left[\begin{array}{cc}
\mathbf{1}_{E_{1}}(\xi) & 0  \tag{6.7}\\
* & \prod_{j=1}^{\infty} m_{2}\left(2^{-j} \xi\right)
\end{array}\right]
$$

where $E_{1}=\bigcap_{j=1}^{\infty} 2^{j}\left(F_{1}\right)=\left[-\frac{4}{7},-\frac{1}{2}\right] \cup\left[-\frac{2}{7}, \frac{2}{7}\right] \cup\left[\frac{1}{2}, \frac{4}{7}\right]$. The lower left entry of the above matrix is represented by a more complicated infinite product which can be computed for some specific choices of the function $v$ satisfying (6.6), see the next example.

The corresponding framelet can be found by choosing a high-pass matrix mask $H$ satisfying (3.6). In addition, if we hope on obtaining a wavelet we should apply Theorem 5.2. Since in our case $\tilde{m}=1$ a.e., we see that $H$ has to be $1 \times 2$ matrix-valued and must satisfy conditions (4.18), (4.19), with the diagonal matrix function $\tilde{\Omega}(\xi)=[1]$. Then, a direct but tedious
calculation shows, that modulo a unimodular $\mathbb{Z}$-periodic function, the high-pass matrix mask $H$ is given by

$$
\left.\begin{array}{rl}
H(\xi)= & {\left[\begin{array}{ll}
\mathbf{1}_{\left.\left[-\frac{1}{4},-\frac{1}{7}\right] \cup \frac{1}{7}, \frac{1}{4}\right]}(\xi) & \mathbf{1}_{\left[-\frac{1}{7},-\frac{1}{14}\right] \cup\left[\frac{1}{14}, \frac{1}{7}\right]}(\xi)
\end{array}\right]} \\
& +e^{2 \pi i \xi} v(\xi+1 / 2)\left[\mathbf{1}_{\left[\frac{3}{7}, \frac{4}{7}\right]}(\xi)\right.
\end{array} \mathbf{1}_{\left[-\frac{1}{14}, \frac{1}{14}\right]}(\xi)\right] \quad \text { for } \xi \in \mathbb{T} .
$$

Define $\psi \in L^{2}(\mathbb{R})$ by $\hat{\psi}=H(\xi) \hat{\Phi}(\xi)$. Then, by Theorem $4.3, \psi$ is a tight framelet for any choice of $\mathbb{Z}$-periodic function $v$ satisfying (6.6).

Note that the choice of $v=\mathbf{1}_{\left[-\frac{1}{14}, \frac{1}{14}\right]}$, corresponds to the matrix mask

$$
M(\xi)=\left[\begin{array}{cc}
\mathbf{1}_{F_{1}}(\xi) & 0 \\
0 & \mathbf{1}_{F_{2}}(\xi)
\end{array}\right],
$$

where $F_{2}$ is the $\mathbb{Z}$-periodization of $\left[-\frac{1}{14}, \frac{1}{14}\right]$. A direct calculation shows that we obtain a tight framelet $\psi$ given by $|\hat{\psi}|=\mathbf{1}_{\left[-\frac{8}{7},-1\right] \cup\left[-\frac{1}{2},-\frac{2}{7}\right] \cup\left[\frac{2}{7}, \frac{1}{2}\right] \cup\left[1, \frac{8}{7}\right]}$. Thus, $\psi$ is not a wavelet. This fact can also be deduced as a consequence of Proposition 5.1. That is, the procedure of constructing refinable vector from low-pass matrix mask can only result in a scaling vector (with the same dimension function) if all but the first column of the product matrix $\prod_{j=1}^{\infty} M\left(2^{-j} \xi\right)$ are zeros. Indeed, if $\psi$ were a wavelet, then by (5.6), condition (5.2) would hold as well. Therefore, the mentioned proposition would imply that (5.1) must be satisfied. However, in our case (5.1) fails. Thus, $\psi$ is a tight framelet, but not a wavelet.

On the other hand, if we choose $v=\mathbf{1}_{\left[\frac{3}{7}, \frac{4}{7}\right]}$, then

$$
M(\xi)=\left[\begin{array}{ll}
\mathbf{1}_{F_{1}}(\xi) & 0 \\
\mathbf{1}_{F_{3}}(\xi) & 0
\end{array}\right]
$$

where $F_{3}$ is the periodization of $\left[\frac{3}{7}, \frac{4}{7}\right]$. A direct calculation shows that we obtain the usual Journé wavelet $\psi$ modified by a negligible unimodular phase factor, i.e., $|\hat{\psi}|=\mathbf{1}_{W}$, see also [3, Example 4.3].

In the next example we construct a large class of non-MSF non-MRA wavelets by an appropriate choice of functions $v$ satisfying (6.6). Naturally, each wavelet in this class must share the dimension function of the Journé wavelet given by (6.1).

Example 6.2. Let $w$ be an arbitrary $\mathbb{Z}$-periodic measurable function satisfying

$$
\begin{equation*}
|w(\xi)|^{2}+|w(\xi+1 / 2)|^{2}=\mathbf{1}_{\left[-\frac{1}{14},-\frac{1}{28}\right] \cup\left[\frac{1}{28}, \frac{1}{14}\right]}(\xi) \quad \text { for a.e. } \xi \in(-1 / 4,1 / 4) \tag{6.8}
\end{equation*}
$$

Then, $v$ given for $\xi \in \mathbb{T}$ by $v(\xi)=w(\xi)+\mathbf{1}_{\left[\frac{13}{28}, \frac{15}{28}\right]}(\xi)$ satisfies (6.6). Define $\mathbb{Z}$-periodic functions

$$
\begin{equation*}
m_{1}(\xi)=w(\xi) \mathbf{1}_{\left[\frac{3}{7}, \frac{13}{28}\right] \cup\left[\frac{15}{28}, \frac{4}{7}\right]}(\xi)+\mathbf{1}_{\left[\frac{13}{28}, \frac{15}{28}\right]}(\xi), \quad m_{2}(\xi)=w(\xi) \mathbf{1}_{\left[-\frac{1}{14},-\frac{1}{28}\right] \cup\left[\frac{1}{28}, \frac{1}{14}\right]}(\xi) . \tag{6.9}
\end{equation*}
$$

Finally, let $M$ be given by (6.2). The same argument as in Example 6.1 shows that $M$ satisfies (2.14)-(2.16) and hence, it is a low-pass matrix mask function. In fact, we obtain a proper
subclass of low-pass matrix masks considered in the previous example. We can also choose a high-pass matrix mask $H$ by emulating Example 6.1. That is, we define

$$
\left.\begin{array}{rll}
H(\xi)= & {\left[\begin{array}{lll}
\mathbf{1}_{\left[-\frac{1}{4},-\frac{1}{7}\right] \cup\left[\frac{1}{7}, \frac{1}{4}\right]}(\xi) & \mathbf{1}_{\left[-\frac{1}{7},-\frac{1}{14}\right] \cup\left[-\frac{1}{28}, \frac{1}{28}\right] \cup\left[\frac{1}{14}, \frac{1}{7}\right]}(\xi)
\end{array}\right]} \\
& +e^{2 \pi i \xi} w(\xi+1 / 2)\left[\mathbf{1}_{\left.\left[\frac{3}{7}, \frac{13}{28}\right] \cup \frac{15}{28}, \frac{4}{7}\right]}(\xi)\right. & \mathbf{1}_{\left[-\frac{1}{14},-\frac{1}{28}\right] \cup\left[\frac{1}{28}, \frac{1}{14}\right]}(\xi)
\end{array}\right] \quad \text { for } \xi \in \mathbb{T} .
$$

The advantage of our choice of low-pass and high-pass matrix masks is twofold. First, the corresponding refinable vector $\Phi$ can be easily computed. Second, it can be shown that $\Phi$ is a scaling vector and the resulting tight framelet $\psi$ is a wavelet. Indeed, note that

$$
M(\xi / 2) M(\xi / 4)=\left[\begin{array}{cc}
\mathbf{1}_{F_{1}}(\xi / 2) \mathbf{1}_{F_{1}}(\xi / 4) & 0  \tag{6.10}\\
m_{1}(\xi / 2) \mathbf{1}_{F_{1}}(\xi / 4)+m_{2}(\xi / 2) m_{1}(\xi / 4) & 0
\end{array}\right]
$$

since $m_{2}(\xi) m_{2}(\xi / 2)=0$ for a.e. $\xi \in \mathbb{R}$. Moreover,

$$
\prod_{j=3}^{\infty} M\left(2^{-j} \xi\right)=\left[\begin{array}{cc}
\mathbf{1}_{E_{1}}(\xi / 4) & 0  \tag{6.11}\\
* & 0
\end{array}\right]
$$

where $E_{1}=\left[-\frac{4}{7},-\frac{1}{2}\right] \cup\left[-\frac{2}{7}, \frac{2}{7}\right] \cup\left[\frac{1}{2}, \frac{4}{7}\right]$. Consequently,

$$
\hat{\Phi}(\xi)=\left[\begin{array}{c}
\hat{\varphi}_{1}(\xi) \\
\hat{\varphi}_{2}(\xi)
\end{array}\right]=\left[\begin{array}{c}
\mathbf{1}_{E_{1}}(\xi) \\
\left(m_{1}(\xi / 2) \mathbf{1}_{F_{1}}(\xi / 4)+m_{2}(\xi / 2) m_{1}(\xi / 4)\right) \mathbf{1}_{E_{1}}(\xi / 4)
\end{array}\right] .
$$

Finally, a direct but tedious calculation shows that

$$
\begin{equation*}
\hat{\varphi}_{2}(\xi)=\mathbf{1}_{\left[-\frac{15}{14},-1\right] \cup\left[1, \frac{15}{14}\right]}+w(\xi / 2) \mathbf{1}_{\left[-\frac{15}{7},-\frac{29}{14}\right] \cup\left[-\frac{8}{7},-\frac{15}{14}\right] \cup\left[\frac{15}{14}, \frac{8}{7}\right] \cup\left[\frac{29}{14}, \frac{15}{7}\right]} . \tag{6.12}
\end{equation*}
$$

To see that $\Phi$ is indeed a scaling vector it suffices to observe that

$$
\sum_{k \in \mathbb{Z}}\left|\hat{\varphi}_{2}(\xi+k)\right|^{2}=\mathbf{1}_{\left[-\frac{1}{7}, \frac{1}{7}\right]}(\xi) \quad \text { for a.e. } \xi \in \mathbb{T},
$$

and check that (2.10) holds.
Finally, one can compute the formula for the corresponding wavelet $\psi=\psi_{w}$,

$$
\begin{align*}
\hat{\psi}(\xi)= & \mathbf{1}_{\left[-\frac{29}{14},-2\right] \cup\left[-\frac{1}{2},-\frac{2}{7}\right] \cup\left[\frac{2}{7}, \frac{1}{2}\right] \cup\left[2, \frac{29}{14}\right]}(\xi)+w(\xi / 4) \mathbf{1}_{\left[-\frac{30}{7},-\frac{29}{7}\right] \cup\left[-\frac{16}{7},-\frac{15}{7}\right] \cup\left[\frac{15}{7}, \frac{16}{7}\right] \cup\left[\frac{29}{7}, \frac{30}{7}\right]}(\xi) \\
& +e^{\pi i \xi} w(\xi / 2+1 / 2) \mathbf{1}_{\left[-\frac{15}{7},-\frac{29}{14}\right] \cup\left[-\frac{8}{7},-\frac{15}{14}\right] \cup\left[\frac{15}{14}, \frac{8}{7}\right] \cup\left[\frac{29}{14}, \frac{, 15}{7}\right]}(\xi) . \tag{6.13}
\end{align*}
$$

Figs. 1 and 2 show graphs of a typical scaling vector $\hat{\Phi}$ and the corresponding wavelet $\hat{\psi}$. We should also add that once the formula (6.13) is established, one can deduce that a wavelet $\psi_{w}$ can be also obtained using interpolation pairs of wavelet sets [19,32].

Observe that the family of wavelets

$$
\mathcal{W}_{\text {nik }}=\left\{\psi_{w}: w \text { satisfies }(6.8)\right\}
$$



Fig. 1. The graphs of $\hat{\varphi}_{1}$ (solid line) and a typical $\hat{\varphi}_{2}$ (dashed line).


Fig. 2. The graph of $\hat{\psi}$ (dashed line corresponds to the part that contains the phase factor $e^{\pi i \xi}$ ).
is pathwise connected in $L^{2}(\mathbb{R})$. Indeed, given two $\mathbb{Z}$-periodic measurable functions $w_{0}$ and $w_{1}$ both satisfying (6.8), it is not difficult to construct a family $\left\{w_{t}\right\}_{t \in[0,1]}$ of functions satisfying (6.8) such that

$$
w_{s}(\xi) \rightarrow w_{t}(\xi) \quad \text { for a.e. } \xi \in \mathbb{T} \text { as } s \rightarrow t
$$

Then, by (6.13), we see that

$$
\hat{\psi}_{w_{s}}(\xi) \rightarrow \hat{\psi}_{w_{t}}(\xi) \quad \text { for a.e. } \xi \in \mathbb{R} \text { as } s \rightarrow t
$$

Since $\left\|\psi_{w_{t}}\right\|=1$ for all $t \in[0,1]$, the map $t \mapsto \psi_{w_{t}}$ is the required continuous path.
Example 6.2 shows that a large class of wavelets can be constructed by the procedure of Theorem 4.3. Moreover, Theorem 5.4 shows that technically every imaginable wavelet can be obtained in that way. However, it is an open problem whether the same is true for all tight framelets.

Two serious difficulties arise when one wants to design a constructive method for obtaining all tight framelets on $\mathbb{R}^{n}$. The first problem is that it is not known if all such framelets are associated to a GMRA. This is often referred to as the "Baggett's problem" [13]. Baggett observed that a
tight framelet $\Psi$ generates a GMRA if and only if its space of negative dilates $V$ satisfies

$$
\begin{equation*}
\bigcap_{j \in \mathbb{Z}} D^{j}(V)=\{0\} . \tag{6.14}
\end{equation*}
$$

We have treated this problem with detail in [16]. An earlier result of the second author [31] assures that if the spectral function of $V$ is integrable, then the above condition is satisfied. It turns out that we can use Lemma 3.1 to improve on this result in the setting of the space of negative dilates. Theorem 6.1 can be also deduced as a consequence of a general result on the intersection of dilates of SI spaces due to the first author [14].

Theorem 6.1. Let $\Psi$ be a tight framelet on $\mathbb{R}^{n}$ with its space of negative dilates $V$. If the set $\left\{\xi \in \mathbb{R}^{n}: \operatorname{dim}_{V}(\xi)<\infty\right\}$ has a positive (Lebesgue) measure, then (6.14) holds and $\Psi$ generates a GMRA.

Proof. As in the proof of Proposition 4.1, let $W=D(V) \ominus V$ and observe that since $\Psi$ consists of a finite number of functions, $W$ has a finite number of generators. That is, we have $\operatorname{dim}_{W} \leqslant N$ for some $N \in \mathbb{N}$. The equation $D(V)=V \oplus W$ implies that

$$
\begin{equation*}
\sum_{d \in \mathcal{D}} m\left(B^{-1} \xi+d\right)=m(\xi)+\operatorname{dim}_{W}(\xi) \leqslant m(\xi)+N \tag{6.15}
\end{equation*}
$$

where $m=\operatorname{dim}_{V}$. Thus, condition (3.11) of Lemma 3.1 is satisfied for such $m$. However, to apply Lemma 3.1 we need to show that $m$ is finite a.e. This can be done using a simple ergodic argument.

Indeed, since the matrix $B: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ preserves the lattice $\mathbb{Z}^{n}$, it induces a measure preserving endomorphism $\tilde{B}: \mathbb{T}^{n} \rightarrow \mathbb{T}^{n}$. Moreover, $\tilde{B}$ is ergodic by [33, Corollary 1.10 .1 ] because $B$ is expansive. Define the set

$$
E=\left\{\xi \in \mathbb{T}^{n}: m(\xi)<\infty\right\} .
$$

The condition (6.15) implies that $\tilde{B}^{-1} E \subset E$. Since $\tilde{B}$ is measure preserving we must have $\tilde{B}^{-1} E=E$ (modulo null sets). Finally, by the ergodicity of $\tilde{B}$, we have either $|E|=0$ or $|E|=1$. Combining this with our hypothesis $|E|>0$, proves that $m(\xi)<\infty$ for a.e. $\xi \in \mathbb{R}^{n}$.

Since all the assumptions of Lemma 3.1 are satisfied for our $m$, we get that $m \in L^{1}\left(\mathbb{T}^{n}\right)$. Equivalently, we have $\sigma_{V} \in L^{1}\left(\mathbb{R}^{n}\right)$. As we mentioned before, the latter implies that (6.14) holds by the result of the second author [31]. Therefore, $\Psi$ generates a GMRA.

If we consider an easier scenario and want to construct all tight framelets associated to a GMRA, we encounter the second difficulty. It is an open problem whether every tight framelet $\Psi$ generating some GMRA $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ can be obtained by the procedure of Theorem 4.3 via the Unitary Extension Principle. In other words, is it possible to find appropriate matrix mask functions $M$ and $H$ resulting by the procedure of Theorem 4.3 in a tight framelet $\Psi$ ? This problem remains open even for tight framelets $\Psi$ associated to an MRA, i.e., when $\operatorname{dim}_{V_{0}} \equiv 1$.

## References

[1] P. Auscher, Solution of two problems on wavelets, J. Geom. Anal. 5 (1995) 181-236.
[2] L. Baggett, An abstract interpretation of the wavelet dimension function using group representations, J. Funct. Anal. 173 (2000) 1-20.
[3] L. Baggett, J. Courter, K. Merrill, The construction of wavelets from generalized conjugate mirror filters in $L^{2}\left(\mathbb{R}^{n}\right)$, Appl. Comput. Harmon. Anal. 13 (2002) 201-223.
[4] L.W. Baggett, P.E.T. Jorgensen, K.D. Merrill, J.A. Packer, Construction of Parseval wavelets from redundant filter systems, J. Math. Phys. 46 (2005) 083502.
[5] L.W. Baggett, P.E.T. Jorgensen, K.D. Merrill, J.A. Packer, A non-MRA $C^{r}$ frame wavelet with rapid decay, Acta Appl. Math. 89 (2005) 251-270.
[6] L. Baggett, H. Medina, K. Merrill, Generalized multi-resolution analyses and a construction procedure for all wavelet sets in $\mathbb{R}^{n}$, J. Fourier Anal. Appl. 5 (1999) 563-573.
[7] L. Baggett, K. Merrill, Abstract harmonic analysis and wavelets in $\mathbf{R}^{n}$, in: The Functional and Harmonic Analysis of Wavelets and Frames, San Antonio, TX, 1999, in: Contemp. Math., vol. 247, Amer. Math. Soc., Providence, RI, 1999, pp. 17-27.
[8] J. Benedetto, S. Li, The theory of multiresolution analysis frames and applications to filter banks, Appl. Comput. Harmon. Anal. 5 (1998) 389-427.
[9] J. Benedetto, O. Treiber, Wavelet frames: Multiresolution analysis and extension principles, in: Wavelet Transforms and Time-Frequency Signal Analysis, in: Appl. Numer. Harmon. Anal., Birkhäuser Boston, Boston, MA, 2001, pp. 3-36.
[10] M. Bownik, The structure of shift-invariant subspaces of $L^{2}\left(\mathbb{R}^{n}\right)$, J. Funct. Anal. 177 (2000) 282-309.
[11] M. Bownik, A characterization of affine dual frames in $L^{2}\left(\mathbf{R}^{n}\right)$, Appl. Comput. Harmon. Anal. 8 (2000) 203-221.
[12] M. Bownik, Anisotropic Hardy spaces and wavelets, Mem. Amer. Math. Soc. 164 (2003), No. 781, 122 pp.
[13] M. Bownik, Baggett's problem for frame wavelets, in: Representations, Wavelets and Frames: A Celebration of the Mathematical Work of Lawrence Baggett, Birkhäuser, 2008, pp. 153-173.
[14] M. Bownik, Intersection of dilates of shift-invariant spaces, Proc. Amer. Math. Soc. 137 (2009) 563-572.
[15] M. Bownik, Z. Rzeszotnik, The spectral function of shift-invariant spaces, Michigan Math. J. 51 (2003) 387-414.
[16] M. Bownik, Z. Rzeszotnik, On the existence of multiresolution analysis for framelets, Math. Ann. 332 (2005) 705720.
[17] M. Bownik, Z. Rzeszotnik, D. Speegle, A characterization of dimension functions of wavelets, Appl. Comput. Harmon. Anal. 10 (2001) 71-92.
[18] L. Brandolini, G. Garrigós, Z. Rzeszotnik, G. Weiss, The behaviour at the origin of a class of band-limited wavelets, in: The Functional and Harmonic Analysis of Wavelets and Frames, San Antonio, TX, 1999, in: Contemp. Math., vol. 247, Amer. Math. Soc., Providence, RI, 1999, pp. 75-91.
[19] X. Dai, D. Larson, Wandering vectors for unitary systems and orthogonal wavelets, Mem. Amer. Math. Soc. 134 (1998), No. 640, viii+68 pp.
[20] I. Daubechies, Ten Lectures on Wavelets, SIAM, Philadelphia, PA, 1992.
[21] I. Daubechies, B. Han, A. Ron, Z. Shen, Framelets: MRA-based constructions of wavelet frames, Appl. Comput. Harmon. Anal. 14 (2003) 1-46.
[22] C. de Boor, R.A. DeVore, A. Ron, The structure of finitely generated shift-invariant spaces in $L_{2}\left(\mathbb{R}^{d}\right)$, J. Funct. Anal. 119 (1994) 37-78.
[23] C. de Boor, R.A. DeVore, A. Ron, Approximation orders of FSI spaces in $L_{2}\left(\mathbb{R}^{d}\right)$, Constr. Approx. 14 (1998) 631-652.
[24] M. Frazier, G. Garrigós, K. Wang, G. Weiss, A characterization of functions that generate wavelet and related expansion, J. Fourier Anal. Appl. 3 (1997) 883-906.
[25] E. Hernández, G. Weiss, A First Course on Wavelets, CRC Press, Boca Raton, FL, 1996.
[26] W. Lawton, Tight frames of compactly supported affine wavelets, J. Math. Phys. 31 (1990) 1898-1901.
[27] M. Paluszyński, H. Šikić, G. Weiss, S. Xiao, Generalized low pass filters and MRA frame wavelets, J. Geom. Anal. 11 (2001) 311-342.
[28] M. Paluszyński, H. Šikić, G. Weiss, S. Xiao, Tight frame wavelets, their dimension functions, MRA tight frame wavelets and connectivity properties, Adv. Comput. Math. 18 (2003) 297-327.
[29] A. Ron, Z. Shen, Affine systems in $L_{2}\left(\mathbb{R}^{d}\right)$ : The analysis of the analysis operator, J. Funct. Anal. 148 (1997) 408-447.
[30] A. Ron, Z. Shen, The wavelet dimension function is the trace function of a shift-invariant system, Proc. Amer. Math. Soc. 131 (2003) 1385-1398.
[31] Z. Rzeszotnik, Calderón's condition and wavelets, Collect. Math. 52 (2001) 181-191.
[32] Z. Rzeszotnik, D. Speegle, On wavelets interpolated from a pair of wavelet sets, Proc. Amer. Math. Soc. 130 (2002) 2921-2930.
[33] P. Walters, An Introduction to Ergodic Theory, Springer-Verlag, New York, 1982.
[34] The Wutam Consortium, Basic properties of wavelets, J. Fourier Anal. Appl. 4 (1998) 575-594.


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